

# Local controllability and non controllability for a 1D wave equation with bilinear control

Karine BEAUCHARD<sup>\*†</sup>

## Abstract

We consider a linear wave equation, on a bounded interval, with bilinear control and Neumann boundary conditions. We study the controllability of this nonlinear control system, locally around a constant reference trajectory. We prove that the following results hold generically.

- For every  $T > 2$ , this system is locally controllable in  $H^3 \times H^2$ , in time  $T$ , with controls in  $L^2((0, T), \mathbb{R})$ .
- For  $T = 2$ , this system is locally controllable up to codimension one in  $H^3 \times H^2$ , in time  $T$ , with controls in  $L^2((0, T), \mathbb{R})$ : the reachable set is (locally) a non flat submanifold of  $H^3 \times H^2$  with codimension one.
- For every  $T < 2$ , this system is not locally controllable, more precisely, the reachable set, with controls in  $L^2((0, T), \mathbb{R})$ , is contained in a non flat submanifold of  $H^3 \times H^2$ , with infinite codimension.

The proof of these results relies on the inverse mapping theorem and second order expansions.

## 1 Introduction

### 1.1 Main result

The goal of this article is to investigate the exact controllability of the wave equation with bilinear controls. We consider the following 1D-wave equation

$$\begin{cases} \frac{\partial^2 w}{\partial t^2}(t, x) = \frac{\partial^2 w}{\partial x^2}(t, x) + u(t)\mu(x)w(t, x), & x \in (0, 1), t \in (0, T), \\ \frac{\partial w}{\partial x}(t, 0) = \frac{\partial w}{\partial x}(t, 1) = 0, \end{cases} \quad (1)$$

where  $\mu \in W^{2,\infty}((0, 1), \mathbb{R})$ . The system (1) is a bilinear control system, in which

- the state is  $(w, \frac{\partial w}{\partial t})$ ,
- the control is the real valued function  $u : [0, T] \rightarrow \mathbb{R}$ .

Let us introduce some conventions and notations. Unless otherwise specified, the functions are real valued. The operator  $A$  is defined by

$$D(A) := \{\varphi \in H^2(0, 1); \varphi'(0) = \varphi'(1) = 0\}, \quad A\varphi := -\varphi''. \quad (2)$$

Its eigenvalues  $(\lambda_k)_{k \in \mathbb{N}}$  and eigenvectors  $(\varphi_k)_{k \in \mathbb{N}}$  are

$$\begin{aligned} \lambda_0 &:= 0, & \varphi_0(x) &:= 1, \\ \lambda_k &:= (k\pi)^2, & \varphi_k(x) &:= \sqrt{2} \cos(k\pi x), \forall k \in \mathbb{N}^*. \end{aligned} \quad (3)$$

---

<sup>\*</sup>CMLA, ENS Cachan, CNRS, Universud, 61 avenue du Président Wilson, F-94230 Cachan, France, email: Karine.Beauchard@cmla.ens-cachan.fr

<sup>†</sup>The author was partially supported by the ‘‘Agence Nationale de la Recherche’’ (ANR), Projet Blanc C-QUID number BLAN-3-139579

We define the spaces

$$H_{(0)}^s(0, 1) := D(A^{s/2}), \forall s > 0 \quad (4)$$

equipped with the norm

$$\|\varphi\|_{H_{(0)}^s} := \left( \sum_{k=0}^{\infty} |k_*^s \langle \varphi, \varphi_k \rangle|^2 \right)^{1/2},$$

where  $k_* := \max\{k, 1\}$ ,  $\forall k \in \mathbb{N}$  and  $\langle \cdot, \cdot \rangle$  is the  $L^2(0, 1)$ -scalar product. Notice that

$$\begin{aligned} H_{(0)}^1(0, 1) &= H^1(0, 1), \\ H_{(0)}^2(0, 1) &= \{\varphi \in H^2(0, 1); \varphi'(0) = \varphi'(1) = 0\}, \\ H_{(0)}^3(0, 1) &= \{\varphi \in H^3(0, 1); \varphi'(0) = \varphi'(1) = 0\}. \end{aligned}$$

The goal of this article is to prove that, under generic assumptions on  $\mu$ , the system (1) is locally controllable around the reference trajectory  $(w(t, x) = 1, u(t) = 0)$ , if and only if  $T > 2$ . The restriction  $T > 2$  is not surprising because this wave equation has a propagation speed equal to 1, but, in this article, a particular attention is given to the case  $T \leq 2$ . Precisely, we prove the following results.

- When  $T > 2$ , the system (1) is locally controllable in  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$  with  $L^2(0, T)$ -controls.
- When  $T = 2$ , the system (1) is not locally controllable in  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$  with  $L^2(0, T)$ -controls because the reachable set is (locally) a non flat submanifold of  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$  with codimension one. However, the system (1) is locally controllable up to codimension one: one can control the couple  $(w - \int_0^1 w(x)dx, \partial w / \partial t)$ . Moreover, for any reachable (local) target, there exists a unique (small) control allowing to reach this target.
- When  $T < 2$ , the system (1) is strongly not controllable: the reachable set, with  $L^2(0, T)$ -controls, is (locally) contained in a non flat submanifold of  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$  with infinite codimension.

The goal of this article is the proof of the following Theorem.

**Theorem 1** *Let  $\mu \in W^{2,\infty}(0, 1)$ . We assume*

$$\exists c > 0 \text{ such that } \frac{c}{k_*^2} \leq |\langle \mu, \varphi_k \rangle|, \forall k \in \mathbb{N}. \quad (5)$$

**(1)** *Let  $T > 2$ . There exists  $\delta > 0$  and a  $C^1$ -map*

$$\begin{aligned} \Gamma_T : \quad \mathcal{V}_T &\rightarrow L^2(0, T) \\ (w_f, \dot{w}_f) &\mapsto \Gamma_T(w_f, \dot{w}_f) \end{aligned}$$

where

$$\mathcal{V}_T := \{(w_f, \dot{w}_f) \in H_{(0)}^3 \times H_{(0)}^2(0, 1); \|w_f - 1\|_{H_{(0)}^3} + \|\dot{w}_f\|_{H_{(0)}^2} < \delta\},$$

such that,  $\Gamma_T(1, 0) = 0$  and for every  $(w_f, \dot{w}_f) \in \mathcal{V}_T$ , the solution of (1) with initial condition

$$\left( w, \frac{\partial w}{\partial t} \right) (0, x) = (1, 0), \forall x \in (0, 1), \quad (6)$$

and control  $u = \Gamma_T(w_f, \dot{w}_f)$  satisfies  $(w, \frac{\partial w}{\partial t})(T) = (w_f, \dot{w}_f)$ .

(2) Let  $T = 2$ . There exists  $\delta, r > 0$  and a  $C^1$  map

$$\begin{aligned} \Gamma_T : \quad \mathcal{V}_T &\rightarrow B_r[L^2(0, T)] \\ (\tilde{w}_f, \dot{w}_f) &\mapsto \Gamma_T(\tilde{w}_f, \dot{w}_f) \end{aligned}$$

where

$$\mathcal{V}_T := \{(\tilde{w}_f, \dot{w}_f) \in H_{(0)}^3 \times H_{(0)}^2(0, 1); \int_0^1 \tilde{w}_f(x) dx = 0, \|\tilde{w}_f\|_{H_{(0)}^3} + \|\dot{w}_f\|_{H_{(0)}^2} < \delta\},$$

$$B_r[L^2(0, T)] := \{u \in L^2((0, T), \mathbb{R}); \|u\|_{L^2} < r\},$$

such that,  $\Gamma_T(0, 0) = 0$  and for every  $(\tilde{w}_f, \dot{w}_f) \in \mathcal{V}_T$ ,  $u \in B_r[L^2(0, T)]$ , the solution of (1), (6) satisfies

$$w(T) - \int_0^1 w(T, x) dx = \tilde{w}_f \text{ and } \frac{\partial w}{\partial t}(T) = \dot{w}_f,$$

if and only if  $u = \Gamma_T(\tilde{w}_f, \dot{w}_f)$ .

The reachable set from (6) is, locally, a  $C^1$ -submanifold with codimension one. More precisely, there exists  $r' > 0$  and a locally surjective non linear  $C^1$ -map  $G_T : H_{(0)}^3 \times H_{(0)}^2(0, 1) \rightarrow \mathbb{R}$  such that, for every  $u \in B_{r'}[L^2(0, T)]$ , the solution of (1), (6) satisfies

$$G_T \left[ \left( w, \frac{\partial w}{\partial t} \right) (T) \right] = 0.$$

(3) We assume

$$\frac{(\mu^2)'(1) \pm (\mu^2)'(0)}{\mu'(1) \pm \mu'(0)} \neq \frac{\int_0^1 \mu(x)^2 dx}{\int_0^1 \mu(x) dx}. \quad (7)$$

Let  $T < 2$ . The reachable set from (6) is, locally, contained in a  $C^1$ -submanifold of  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$ , with infinite codimension, that does not coincide with its tangent space at  $(1, 0)$ . More precisely, there exists  $r > 0$ , a strict vector subspace  $R_T$  of  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$  with infinite dimension and a locally surjective  $C^1$  map

$$G_T : H_{(0)}^3 \times H_{(0)}^2(0, 1) \rightarrow R_T$$

such that, for every  $u \in B_r[L^2(0, T)]$ , the solution of (1), (6) satisfies

$$G_T \left[ \left( w, \frac{\partial w}{\partial t} \right) (T) \right] = 0.$$

**Remark 1** Notice that, when (5) holds, then  $\int_0^1 \mu = \langle \mu, \varphi_0 \rangle \neq 0$  and  $\mu'(1) \pm \mu'(0) \neq 0$ . Indeed, we have

$$\langle \mu, \varphi_k \rangle = \frac{\sqrt{2}}{(k\pi)^2} \left( (-1)^k \mu'(1) - \mu'(0) \right) - \frac{\sqrt{2}}{(k\pi)^2} \int_0^1 \mu''(x) \cos(k\pi x) dx. \quad (8)$$

This remark gives a sense to each term in (7).

**Remark 2** The assumptions (5) and (7) hold simultaneously, for example, with  $\mu(x) = x^2$ , because

$$\begin{aligned} \langle x^2, \varphi_0 \rangle &= \int_0^1 x^2 dx = \frac{1}{3}, \\ \langle x^2, \varphi_k \rangle &= \int_0^1 x^2 \sqrt{2} \cos(k\pi x) dx = \frac{(-1)^k 2\sqrt{2}}{(k\pi)^2}, \forall k \in \mathbb{N}^*, \end{aligned} \quad (9)$$

$$\frac{(\mu^2)'(1) \pm (\mu^2)'(0)}{\mu'(1) \pm \mu'(0)} = 2, \text{ and } \frac{\int_0^1 \mu(x)^2 dx}{\int_0^1 \mu(x) dx} = \frac{3}{5}.$$

But (5) and (7) are not always satisfied. For example, (5) does not hold when  $\langle \mu, \varphi_k \rangle = 0$  for some  $k \in \mathbb{N}$ , or when  $\mu$  has a symmetry with respect to  $x = 1/2$ . However, the assumptions (5) and (7) are generic in  $W^{2,\infty}(0,1)$  (see Appendix A for a proof), thus, Theorem 1 is very general.

**Remark 3** In Theorem 1, the spaces are optimal. Indeed, we will see in this article that, for every control  $u \in L^2(0,T)$ , there exists a unique solution of (1), (6) and it satisfies

$$\left( w, \frac{\partial w}{\partial t} \right) (T) \in H_{(0)}^3 \times H_{(0)}^2(0,1).$$

**Remark 4** Let us mention the reference [15] by Coron, Rouchon and the author, in which similar results are proved. In this reference, we consider the Bloch equation

$$\frac{\partial M}{\partial t}(t, \omega) = \begin{pmatrix} 0 & -\omega & v(t) \\ \omega & 0 & u(t) \\ -v(t) & -u(t) & 0 \end{pmatrix} M(t, \omega), t \in [0, +\infty), \omega \in (\omega_*, \omega^*),$$

where  $-\infty \leq \omega_* < \omega^* \leq +\infty$ ,  $u, v : [0, +\infty) \rightarrow \mathbb{R}$ . It is a control system where the state is the function  $M = M(t, \omega)$  and the control is  $(u, v) : [0, +\infty) \rightarrow \mathbb{R}^2$ . This system is a prototype for infinite dimensional bilinear control systems, with continuous spectrum. In [15, Theorem 2], we prove that, when  $\omega_* = -\infty$  and  $\omega^* = +\infty$ , then, this system is not exactly controllable, locally around the reference trajectory  $(M_{ref} = e_3, u_{ref} = 0, v_{ref} = 0)$ . The proof consists in proving that the reachable set from  $M(0, \omega) = e_3$ , in time  $T$ , with bounded  $L^2(0, T)$ -controls, is locally a non flat submanifold of some functional space, with infinite codimension. The proof of this result relies on the inverse mapping theorem, and second order expansions, as in the present article.

In this article, the same letter  $C$  denotes a positive constant that can change from one line to another one.

## 1.2 Sketch of the proof

The proof of Theorem 1 relies on the inverse mapping theorem, applied to the end point map

$$\Theta_T : u \mapsto \left( w, \frac{\partial w}{\partial t} \right) (T), \quad (10)$$

where  $w$  solves (1), (6).

First, we prove that, for every  $T > 0$ , the map  $\Theta_T$  is  $C^1$  between the following spaces

$$\Theta_T : L^2(0, T) \rightarrow H_{(0)}^3 \times H_{(0)}^2(0, 1).$$

Then, the local controllability of the nonlinear system when  $T > 2$  (i.e. the local surjectivity of  $\Theta_T$ ) is a consequence of the surjectivity of  $d\Theta_T(0)$ . And the non controllability of the nonlinear system when  $T \leq 2$  is a consequence of the injectivity and non-surjectivity of  $d\Theta_T(0)$ . More precisely, we prove the following results.

- When  $T > 2$ , the continuous linear map  $d\Theta_T(0) : L^2(0, T) \rightarrow H_{(0)}^3 \times H_{(0)}^2(0, 1)$  has a continuous right inverse. This means that the linearized system around the reference trajectory  $(w(t, x) = 1, u(t) = 0)$  is controllable, in time  $T$ , in  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$ , with controls in  $L^2(0, T)$ .

- When  $T = 2$ , the continuous linear map  $d\Theta_T(0) : L^2(0, T) \rightarrow H_{(0)}^3 \times H_{(0)}^2(0, 1)$  is injective, its image  $R_T$  is a vector subspace of  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$  with codimension one, and the map  $d\Theta_T(0) : L^2(0, T) \rightarrow R_T$  has a continuous (left and right) inverse. This means that the linearized system around the reference trajectory ( $w(t, x) = 1, u(t) = 0$ ) is controllable up to codimension one, in time  $T$ , in  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$ , with controls in  $L^2(0, T)$ : it misses exactly one direction. Moreover, for every reachable target, there exists a unique control allowing this motion.
- When  $T < 2$ , the continuous linear map  $d\Theta_T(0) : L^2(0, T) \rightarrow H_{(0)}^3 \times H_{(0)}^2(0, 1)$  is injective, its image  $R_T$  is a vector subspace of  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$  with infinite codimension and the map  $d\Theta_T(0) : L^2(0, T) \rightarrow R_T$  has a continuous (left and right) inverse. This means that the linearized system around the reference trajectory ( $w(t, x) = 1, u(t) = 0$ ) is strongly not controllable: it misses an infinite number of directions. Moreover, for every reachable target, there exists a unique control allowing this motion.

Thus, by applying the inverse mapping theorem, we prove that the reachable set in time  $T \leq 2$  is a strict submanifold of  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$ . Now, let us explain how we prove this submanifold is not flat. First, we prove that the image of the quadratic form  $d^2\Theta_T(0)$  is not contained in the image of the linear map  $d\Theta_T(0)$ . Then, thanks to a second order expansion of  $\Theta_T$  around 0, we see that the (local) submanifold (i.e. the image of  $\Theta_T$ ) does not coincide with its tangent space at  $(1, 0)$  (i.e. the image of  $d\Theta_T(0)$ ).

**Remark 5** *The first (local) exact controllability result, for an infinite dimensional bilinear system, has been proved in [10], for a Schrödinger equation. In [10], the strategy is the same as in this article: first, we prove the controllability of the linearized system and then, we conclude by applying an inverse mapping theorem. However, because of an a priori loss of regularity, we use the Nash-Moser implicit function theorem, instead of the classical inverse mapping theorem. Thus, the analysis is quite complicated.*

*One of the interests of this article is to provide an example of infinite dimensional bilinear control system (i.e. the equation (1)), for which the proof of the (local) exact controllability relies only on the classical inverse mapping theorem, and is rather simple. In order to avoid the use of the Nash-Moser theorem, we emphasize a 'hidden' regularization effect for the equation 1.*

## 1.3 A brief bibliography

### 1.3.1 A previous negative result for this equation

The following result is due to Ball, Marsden and Slemrod [5, Theorem 3.6].

**Theorem 2** *Let  $X$  be a Banach space with infinite dimension. Let  $\mathcal{A}$  be the generator of a  $C^0$ -group of bounded operators of  $X$  and  $\mathcal{B}$  be a bounded operator of  $X$ . For  $w_0 \in X$  and  $p \in L_{loc}^1([0, +\infty), \mathbb{R})$ ,  $U[T; p, w_0]$  denotes the value at time  $T$  of the unique weak solution of*

$$\begin{cases} \frac{dw}{dt} = \mathcal{A}w + p(t)\mathcal{B}w(t), \\ w(0) = w_0. \end{cases} \quad (11)$$

*For every  $w_0 \in X$ , the reachable set from  $w_0$ ,*

$$\mathcal{R}(w_0) := \{U[T; p, w_0]; T \geq 0, p \in L_{loc}^r([0, +\infty), \mathbb{R}), r > 1\}$$

*has an empty interior in  $X$ .*

A consequence of this theorem is the non controllability of the system (11), in  $X$ , with controls  $p \in L_{loc}^r([0, +\infty), \mathbb{R})$ ,  $r > 1$ .

Theorem 2 applies to the system (1), written in first order form, with

$$\begin{aligned} X &:= H_{(0)}^2 \times H^1(0, 1), \\ D(\mathcal{A}) &:= H_{(0)}^2 \times H^1(0, 1), \quad \mathcal{A} := \begin{pmatrix} 0 & I \\ -A & 0 \end{pmatrix}, \\ D(\mathcal{B}) &:= L^2 \times L^2(0, 1), \quad \mathcal{B} := \begin{pmatrix} 0 & 0 \\ \mu & 0 \end{pmatrix}. \end{aligned} \tag{12}$$

Indeed, for every  $(w_0, \dot{w}_0) \in H^1 \times L^2(0, 1)$ , we have

$$e^{\mathcal{A}t} \begin{pmatrix} w_0 \\ \dot{w}_0 \end{pmatrix} = \begin{pmatrix} w(t) \\ \dot{w}(t) \end{pmatrix},$$

where

$$\begin{aligned} w(t) &= (\langle w_0, \varphi_0 \rangle + \langle \dot{w}_0, \varphi_0 \rangle t) \varphi_0 + \sum_{k=1}^{\infty} \left( \langle w_0, \varphi_k \rangle \cos(\sqrt{\lambda_k} t) + \frac{1}{\sqrt{\lambda_k}} \langle \dot{w}_0, \varphi_k \rangle \sin(\sqrt{\lambda_k} t) \right) \varphi_k, \\ \dot{w}(t) &= \langle \dot{w}_0, \varphi_0 \rangle \varphi_0 + \sum_{k=1}^{\infty} \left( -\sqrt{\lambda_k} \langle w_0, \varphi_k \rangle \sin(\sqrt{\lambda_k} t) + \langle \dot{w}_0, \varphi_k \rangle \cos(\sqrt{\lambda_k} t) \right) \varphi_k. \end{aligned}$$

Thus  $\mathcal{A}$  generates a  $C^0$ -group of bounded operators of  $X$ . Moreover  $\mathcal{B}$  is a bounded operator of  $X$  when  $\mu \in W^{1,\infty}(0, 1)$ . For a precise definition of weak solutions of (1), we refer to Proposition 2. Thanks to Theorem 2, we have the following non controllability result for (1).

**Proposition 1** *Let  $\mu \in W^{1,\infty}(0, 1)$ ,  $T > 0$  and  $(w_0, \dot{w}_0) \in H_{(0)}^2 \times H^1(0, 1)$ . For  $u \in L_{loc}^1[0, +\infty)$ ,  $U[T; u, w_0, \dot{w}_0]$  denotes the value at time  $T$  of the weak solution of (1) with initial condition*

$$\left( w, \frac{\partial w}{\partial t} \right) (0) = (w_0, \dot{w}_0).$$

*The reachable set from  $(w_0, \dot{w}_0)$ ,*

$$\mathcal{R}(w_0, \dot{w}_0) := \{U[T; u, w_0, \dot{w}_0]; T > 0, u \in L_{loc}^r[0, +\infty), r > 1\}$$

*has an empty interior in  $H_{(0)}^2 \times H^1(0, 1)$ .*

Thus, the system (1) is not controllable in  $H_{(0)}^2 \times H^1(0, 1)$  with controls in  $L_{loc}^r[0, +\infty)$ ,  $r > 1$ .

**Remark 6** *Notice that Theorem 2 does not apply with*

$$X := H_{(0)}^3 \times H_{(0)}^2(0, 1).$$

*Indeed,  $\mathcal{A}$  generates a  $C^0$ -group of bounded operators of  $X$ , but  $\mathcal{B}$  does not map  $X$  into  $X$ : for  $\varphi \in H_{(0)}^3(0, 1)$  (i.e.  $\varphi \in H^3(0, 1)$  and  $\varphi'(0) = \varphi'(1) = 0$ ), we have  $(\mu\varphi)'(0) = \mu'(0)\varphi(0)$  and  $(\mu\varphi)'(1) = \mu'(1)\varphi(1)$  that may not vanish.*

Such a negative controllability result may be rather weak, because it does not prevent from positive controllability results, in different functional spaces. For example, the reachable set  $\mathcal{R}(w_0, \dot{w}_0)$  may be the whole space  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$  (which has an empty interior in  $H_{(0)}^2 \times H_{(0)}^1(0, 1)$ ) and then the system would be controllable in  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$ . In this article, we prove that this is indeed the case, at least locally, when  $T > 2$ . On the contrary, when  $T < 2$ , the system (1) is not controllable in a very strong sense (stronger than Ball,

Marsden and Slemrod's one): the reachable set  $\mathcal{R}(1, 0)$  is locally a non flat submanifold of  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$ , with infinite codimension. In particular, when  $T < 2$ , no positive exact controllability result can be expected in smoother spaces (because the manifold is not flat). Thus, the results of this article complete the ones of [5].

The same kind of situation arises with bilinear Schrödinger or beam equations (see [13], [10], [11], [12], [14]).

### 1.3.2 Iterated Lie brackets for general bilinear systems

Now, let us discuss the exact controllability of general bilinear systems.

First, the controllability of *finite dimensional* bilinear control systems (i.e. modeled by an ordinary differential equation) is well understood. Let us consider the control system

$$\frac{dX}{dt} = AX + u(t)BX, \quad (13)$$

where  $X(t) \in \mathbb{R}^n$  is the state,  $A, B$  are  $n \times n$  matrices, and  $t \mapsto u(t) \in \mathbb{R}$  is the control. The controllability of (13) is linked to the rank of the Lie algebra spanned by  $A$  and  $B$  (see for example [2] by Agrachev and Sachkov, [21, Chapter 3] by Coron or [22] by D'Alessandro).

In *infinite dimension*, there are cases where the iterated Lie brackets provide the right intuition. For instance, it holds for the non controllability of the harmonic quantum oscillator with bilinear control (see [37] by Mirrahimi and Rouchon). However, the Lie brackets are sometimes less powerful in infinite dimension than in finite dimension. It is precisely the case of our system. Let us compute the iterated Lie brackets of the operators  $\mathcal{A}$  and  $\mathcal{B}$  defined by (12), at the point

$$\mathcal{W}_0 = \begin{pmatrix} 1 \\ 0 \end{pmatrix}.$$

We have

$$[\mathcal{A}, \mathcal{B}]\mathcal{W}_0 = (\mathcal{A}\mathcal{B} - \mathcal{B}\mathcal{A})\mathcal{W}_0 = \mathcal{A}\mathcal{B}\mathcal{W}_0 = \begin{pmatrix} \mu \\ 0 \end{pmatrix}$$

( $\mu$  is assumed to belong to  $W^{2,\infty}(0, 1)$ ). Notice that  $[\mathcal{A}, \mathcal{B}]\mathcal{W}_0$  does not belong to  $D(\mathcal{A})$  because  $\mu'$  may not vanish at 0 and 1. Thus, in order to compute the iterated Lie bracket  $[\mathcal{A}, [\mathcal{A}, \mathcal{B}]]\mathcal{W}_0$ , one needs to extend the definition of  $\mathcal{A}$  to couples  $(w_0, w_1)$  such that  $w'_0$  does not vanish at 0 and 1. A natural choice is

$$\mathcal{A} \begin{pmatrix} w_0 \\ w_1 \end{pmatrix} := \begin{pmatrix} w_1 \\ w''_0 - w'_0(1)\delta_1 + w'_0(0)\delta_0 \end{pmatrix}, \forall (w_0, w_1) \in H^2 \times H^1(0, 1). \quad (14)$$

With this definition, we get

$$\begin{aligned} [\mathcal{A}, [\mathcal{A}, \mathcal{B}]]\mathcal{W}_0 &= \begin{pmatrix} 0 \\ \mu'' - \mu'(1)\delta_1 + \mu'(0)\delta_0 \end{pmatrix}, \\ [\mathcal{A}, [\mathcal{A}, [\mathcal{A}, \mathcal{B}]]]\mathcal{W}_0 &= \begin{pmatrix} \mu'' - \mu'(1)\delta_1 + \mu'(0)\delta_0 \\ 0 \end{pmatrix}. \end{aligned}$$

But again,  $[\mathcal{A}, [\mathcal{A}, [\mathcal{A}, \mathcal{B}]]]\mathcal{W}_0$  does not belong to  $H^2 \times H^1(0, 1)$ , thus the definition (14) cannot be used to compute  $[\mathcal{A}, [\mathcal{A}, [\mathcal{A}, [\mathcal{A}, \mathcal{B}]]]]\mathcal{W}_0$ . Moreover, even if we could give a sense to any iterated Lie bracket, because of the presence of Dirac masses, it would not be clear which space the Lie brackets should generate in case of local controllability around the references trajectory ( $w(t, x) = 1, u(t) = 0$ ). Therefore, the way the Lie algebra rank condition

could be used directly in infinite dimension is not clear.

Finally, let us quote important articles about the controllability of PDEs, in which positive results are proved by applying such geometric control methods but to the (finite dimensional) Galerkin approximations of the equation. In [3] by Sarychev and Agrachev and [41] by Shirikyan, the authors prove exact controllability results for dissipative equations. In [19] by Boscain, Chambrion, Mason and Sigalotti, the authors prove approximate controllability results for Schrödinger equations.

### 1.3.3 Wave equation with bilinear control

Now, let us quote few articles about the controllability of wave equations with bilinear control. In [32], Khapalov considers the following control system

$$\begin{cases} \frac{\partial^2 y}{\partial t^2} = \frac{\partial^2 y}{\partial x^2} + v(t, x)y(t, x) - \gamma(t)y_t - F(t, x, y), x \in (0, 1), t \in (0, +\infty), \\ y(t, 0) = y(t, 1) = 0, \end{cases} \quad (15)$$

in which the controls are  $v \in L^\infty((0, +\infty) \times (0, 1))$  and  $\gamma \in L^\infty(0, +\infty)$ . This equation represents a semilinear vibrating string, with clamped ends, with a variable axial load  $v(t, x)$  and a variable damping gain  $\gamma(t)$ . The nonlinearity  $F$  is fixed. Such controllability problems may arise in the context of 'smart materials', whose properties can be altered by applying various factors (temperature, electric current, magnetic field). In [32], the author proves the global approximate controllability to nonnegative equilibrium states:  $\forall (y_0, y_1) \in H_0^1 \times L^2(0, 1)$  with  $(y_0, y_1) \neq 0$ ,  $\forall y_d \in L^1(0, 1)$  with  $y_d \geq 0$  a.e. on  $(0, 1)$ ,  $\forall \epsilon > 0$  there exists  $T = T(\epsilon, y_0, y_1, y_d) > 0$  and piecewise-constant-in-time controls  $(v, \gamma)$  such that the solution of (15) with initial condition

$$\left( y, \frac{\partial y}{\partial t} \right) (0) = (y_0, y_1),$$

satisfies

$$\|y(T) - y_d\|_{L^2} + \left\| \frac{\partial y}{\partial t} \right\|_{L^2} < \epsilon.$$

The proof consists in, first, finding a control  $(v, \gamma)$  that realizes the approximate controllability for the homogeneous truncated system (i.e. with  $F = 0$ ), and then, proving that, the nonlinear system with the same control follows closely the linear one. We also refer to [31] and [29] by Khapalov for similar results on similar equations (with  $\gamma = 0$  or  $F = 0$ ).

### 1.3.4 Wave equation with linear controls

Now, let us quote few articles about the controllability of wave equations with distributed or boundary controls acting linearly on the state. There is a huge literature on this subject. One of the best result has been obtained by Bardos, Lebeau and Rauch in [6]. See also the paper [18] by Burq and Gérard, the paper [17] by Burq for improvements or simpler proofs, and the papers [43] by Zuazua for semilinear equations. Let us also mention the survey paper [40] by Russell and the books [21] by Coron, [25] by Fursikov and Imanuvilov, [35] by Jacques Louis Lions and [34] by Komornik, where one can find plenty of results and useful references.

### 1.3.5 Other results about infinite dimensional bilinear systems

In the last years, important progress have been made about the controllability of Schrödinger equations with bilinear control.

The first results were negative: in [42], Turinici adapted Theorem 2 to linear Schrödinger equations; in [28], Lange and Teismann adapted it to nonlinear equations; in [37], Mirrahimi and Rouchon proved a stronger negative result for the quantum harmonic oscillator.



Concerning exact controllability issues, local results for 1D models have been proved in [10, 11] by the author, who proposed a simplified proof in [13]; almost global results have been proved in [14], by Coron and the author. In [20], Coron proved the existence of a positive minimal time required for the local controllability of the 1D model studied in [10].

Now, let us quote some approximate controllability results. In [16] Mirrahimi and the author proved the global approximate controllability, in infinite time, for a 1D model and in [36] Mirrahimi proved a similar result for equations involving a continuous spectrum. Approximate controllability, in finite time, has been proved for particular models by Boscain and Adami in [1], by using adiabatic theory and intersection of the eigenvalues in the space of controls. Approximate controllability, in finite time, for more general models, have been studied by 3 teams, with different tools: by Boscain, Chambrion, Mason, Sigalotti in [19], with geometric control methods; by Nersesyan in [39, 38] with feedback controls and variational methods; and by Ervedoza and Puel in [24] thanks to a simplified model.

Let us emphasize that the local exact controllability of [13] and the global approximate controllability of [39, 38] can be put together in order to get the global exact controllability of 1D models (see [38]).

Optimal control techniques have also been investigated for Schrödinger equations with a non linearity of Hartee type in [7, 8] by Baudouin, Kavian, Puel and in [23] by Cances, Le Bris, Pilot. An algorithm for the computation of such optimal controls is studied in [9] by Baudouin and Salomon.

Finally, let us also quote [30, 33] by Khapalov for approximate controllability results about the heat equation, [12] by the author for an exact controllability result about a 1D beam equation, and [15] for a negative exact controllability result and positive approximate controllability results for the Bloch equation.

#### 1.4 A toy model for 2D quantum systems

Finally, let us emphasize that the system (1) may be considered as a toy model for 2D (i.e.  $n = 2$ ) Schrödinger bilinear control systems,

$$\begin{cases} i \frac{\partial \psi}{\partial t} = -\Delta \psi - u(t)\mu(x)\psi, x \in \Omega, t \in [0, T], \\ \psi(t, x) = 0, x \in \partial\Omega, \end{cases} \quad (16)$$

where  $\Omega$  is a bounded regular open subset of  $\mathbb{R}^n$ , and  $\mu : \Omega \rightarrow \mathbb{R}$  is a smooth function.

The system (16) represents a quantum particle in a infinite square potential well  $\Omega$ , subjected to a 1D uniform (in space) time dependent electric field with amplitude  $u(t)$ . The function  $\mu$  is the dipolar moment of the particle. The controllability of such systems is a challenging problem.

In the references above, the approximate controllability results [19, 39, 38] hold in any space dimension ( $\forall n \in \mathbb{N}^*$ ), but the local exact controllability results [13] hold only in 1D ( $n = 1$ ). Thus, the global exact controllability is proved only in 1D (see [38]). It would be interesting to know if the same program works in any dimension, i.e. if the local exact controllability result also holds in 2D and 3D.

A key point in the proof of [13] is the following property: the eigenvalues of the Laplacian on a 1D domain (take, for instance  $\lambda_k = (k\pi)^2$ ,  $k \in \mathbb{N}^*$  with  $\Omega = (0, 1)$ ) satisfy a gap condition:

$$\exists \delta > 0 \text{ such that } \lambda_{k+1} - \lambda_k \geq \delta, \forall k \in \mathbb{N}^*.$$

Such a property does not hold on 2D and 3D domains, for which we only know the Weyl formula,

$$\exists d > 0, \alpha \in (0, n/2) \text{ such that } \text{Card}\{k \in \mathbb{N}; \mu_k \in [0, t]\} = dt^{n/2} + O(t^\alpha) \text{ when } t \rightarrow +\infty. \quad (17)$$

The system (1) may be considered as a toy model for (16) with  $n = 2$ . Indeed, the spectrum of the underlying operator  $\mathcal{A}$  defined by (12) satisfies the Weyl formula (17) with  $n = 2$ , : its eigenvalues are  $(ik\pi)_{k \in \mathbb{N}}$  with the associated eigenvectors  $(X_k)_{k \in \mathbb{N}}$ ,

$$X_k := \begin{pmatrix} \varphi_k \\ ik\pi\varphi_k \end{pmatrix}, \forall k \in \mathbb{N}^*$$

(see (3) for a definition of  $\varphi_k$ ). The control system (1) is easier to deal with than (16) because the spectrum of the underlying operator has more structure.

## 1.5 Structure of this article

This article is organized as follows.

The Section 2 is dedicated to the well posedness of the Cauchy problem (1), (6).

In Subsection 2.1, we state classical results about existence, uniqueness, regularity, and bounds for the solutions of a more general Cauchy problem.

In Subsection 2.2, improving these classical results, we prove that the end point map  $\Theta_T$ , defined by (10), is  $C^1$  from  $L^2(0, T)$  to  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$ .

In Section 3 we consider the linearized system of (1) around the reference trajectory  $(w(t, x) = 1, u(t) = 0)$ . We study its controllability in  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$  with  $L^2(0, T)$ -controls.

In Section 4, we study the second order term around  $(w(t, x) = 1, u(t) = 0)$ . We prove that, for every  $T \geq 2$ , the image of the quadratic form  $d^2\Theta_T(0)$  is not contained in the image of the linear map  $d\Theta_T(0)$ .

In Section 5 we prove Theorem 1, by applying the inverse mapping theorem.

Finally, Section 6 is dedicated to conclusions, open problems and perspectives.

## 2 Well posedness and $C^1$ regularity of the end point map

This section is dedicated to the statement of existence, uniqueness, regularity results, and bounds for the solutions of the Cauchy problem

$$\begin{cases} \frac{\partial^2 w}{\partial t^2}(t, x) = \frac{\partial^2 w}{\partial x^2}(t, x) + u(t)\mu(x)w(t, x) + f(t, x), & x \in (0, 1), t \in \mathbb{R}_+, \\ \frac{\partial w}{\partial x}(t, 0) = \frac{\partial w}{\partial x}(t, 1) = 0, \\ w(0, x) = w_0(x), \\ \frac{\partial w}{\partial t}(0, x) = \dot{w}_0(x). \end{cases} \quad (18)$$

These results are presented in Subsection 2.1. Then, in Subsection 2.2, improving the results of Subsection 2.1, we prove that the map  $\Theta_T$ , defined by (10), is of class  $C^1$  from  $L^2(0, T)$  to  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$ .

### 2.1 Existence, uniqueness, regularity and bounds

In order to study the well posedness of (18), it is convenient to write it in first order form. With the notations

$$\mathcal{W} := \begin{pmatrix} w \\ \frac{\partial w}{\partial t} \end{pmatrix}, \quad \mathcal{W}_0 := \begin{pmatrix} w_0 \\ \dot{w}_0 \end{pmatrix}, \quad \mathcal{F}(t, x) := \begin{pmatrix} 0 \\ f(t, x) \end{pmatrix},$$

and  $\mathcal{A}, \mathcal{B}$  defined by (12), the equation (18) may be written

$$\begin{cases} \frac{\partial \mathcal{W}}{\partial t}(t, x) = \mathcal{A}\mathcal{W}(t, x) + u(t)\mathcal{B}\mathcal{W}(t, x) + \mathcal{F}(t, x), \\ \mathcal{W}(0) = \mathcal{W}_0. \end{cases} \quad (19)$$

The operator  $\mathcal{A}$  generates a  $C^0$ -group of bounded operators of  $H_{(0)}^{s+1} \times H_{(0)}^s(0, 1)$ , for every  $s \geq 0$  (see (4) for a definition) and the operator  $\mathcal{B}$  is bounded on  $H_{(0)}^2 \times H^1(0, 1)$  when  $\mu \in W^{1,\infty}(0, 1)$ . These two facts allow to prove the following classical existence result of weak solutions for (18).

**Proposition 2** *Let  $\mu \in W^{1,\infty}(0, 1)$  and  $T > 0$ . There exists  $C = C(\mu, T) > 0$  such that, for every  $u \in L^1(0, T)$ ,  $(w_0, \dot{w}_0) \in H_{(0)}^2 \times H^1(0, 1)$ , and  $f \in L^1((0, T), H^1(0, 1))$ , there exists a unique weak solution of (18), i.e. a function*

$$\left(w, \frac{\partial w}{\partial t}\right) \in C^0([0, T], H_{(0)}^2 \times H^1(0, 1))$$

such that the following equality holds in  $H_{(0)}^2 \times H^1(0, 1)$ , for every  $t \in [0, T]$ ,

$$\mathcal{W}(t) = e^{\mathcal{A}t}\mathcal{W}_0 + \int_0^t e^{\mathcal{A}(t-\tau)} \left(u(\tau)\mathcal{B}\mathcal{W}(\tau) + \mathcal{F}(\tau)\right) d\tau, \quad (20)$$

and this weak solution satisfies

$$\left\| \left(w, \frac{\partial w}{\partial t}\right) \right\|_{C^0([0, T], H_{(0)}^2 \times H^1)} \leq C \left( \|(w_0, \dot{w}_0)\|_{H_{(0)}^2 \times H^1} + \|f\|_{L^1((0, T), H^1)} \right) e^{C\|u\|_{L^1}}. \quad (21)$$

**Proof of Proposition 2 :** The existence and uniqueness come from a fixed point argument on the map  $F$  defined on  $C^0([0, T], H_{(0)}^2 \times H^1(0, 1))$  by  $F(\mathcal{W}) := \xi$  where

$$\xi(t) = e^{\mathcal{A}t}\mathcal{W}_0 + \int_0^t e^{\mathcal{A}(t-\tau)} \left(u(\tau)\mathcal{B}\mathcal{W}(\tau) + \mathcal{F}(\tau)\right) d\tau, \forall t \in [0, T].$$

$F$  maps  $C^0([0, T], H_{(0)}^2 \times H^1(0, 1))$  into itself because  $\mathcal{B}$  and  $e^{\mathcal{A}t}$  preserve  $H_{(0)}^2 \times H^1(0, 1)$ . When  $\|u\|_{L^1((0, T), \mathbb{R})}$  is small enough, then  $F$  is a contraction, because

$$\begin{aligned} \|F(\mathcal{W}_1)(t) - F(\mathcal{W}_2)(t)\|_{H_{(0)}^2 \times H^1} &= \left\| \int_0^t e^{\mathcal{A}(t-\tau)} u(\tau) \mathcal{B} (\mathcal{W}_1(\tau) - \mathcal{W}_2(\tau)) d\tau \right\|_{H_{(0)}^2 \times H^1} \\ &\leq \int_0^t |u(\tau)| \left\| e^{\mathcal{A}(t-\tau)} \mathcal{B} (\mathcal{W}_1(\tau) - \mathcal{W}_2(\tau)) \right\|_{H_{(0)}^2 \times H^1} d\tau \\ &\leq C_1 \int_0^t |u(\tau)| \left\| \mathcal{B} (\mathcal{W}_1(\tau) - \mathcal{W}_2(\tau)) \right\|_{H_{(0)}^2 \times H^1} d\tau \\ &\leq C_1 C_2 \|u\|_{L^1(0, T)} \|\mathcal{W}_1 - \mathcal{W}_2\|_{C^0([0, T], H_{(0)}^2 \times H^1)}, \end{aligned}$$

where  $C_1 = C_1(\mathcal{A}, T)$ ,  $C_2 = C_2(\mathcal{B}) > 0$ . Thus,  $F$  has a unique fixed point  $\mathcal{W} \in C^0([0, T], H_{(0)}^2 \times H^1)$  that satisfies (20). If  $\|u\|_{L^1((0, T), \mathbb{R})}$  is not small, one may use  $0 = T_0 < T_1 < \dots < T_n = T$  where, for  $i = 0, \dots, n-1$ ,  $\|u\|_{L^1(T_i, T_{i+1})}$  is small enough so that the previous result holds on  $[T_i, T_{i+1}]$ , for  $i = 0, \dots, n-1$ . Then we glue the solutions defined on  $[T_0, T_1]$ ,  $[T_1, T_2]$ , ...,  $[T_{n-1}, T_n]$ . We deduce from the equality (20) that

$$\|\mathcal{W}(t)\|_{H_{(0)}^2 \times H^1} \leq C_1 \left( \|\mathcal{W}_0\|_{H_{(0)}^2 \times H^1} + \|F\|_{L^1((0, T), H_{(0)}^2 \times H^1)} + \int_0^t |u(\tau)| C_2 \|\mathcal{W}(\tau)\|_{H_{(0)}^2 \times H^1} d\tau \right),$$

and Gronwall's Lemma gives (21).  $\square$

**Remark 7** *This proof does not work with  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$  instead of  $H_{(0)}^2 \times H^1(0, 1)$  because  $\mathcal{B}$  does not conserve  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$ . Indeed, for  $\varphi \in H_{(0)}^3(0, 1)$  (i.e.  $\varphi \in H^3(0, 1)$  and  $\varphi'(0) = \varphi'(1) = 0$ ), we have  $(\mu\varphi)' = \mu'\varphi$  at  $x = 0, 1$  that may not vanish. Thus it is not obvious that the map  $\Theta_T$  defined by (10) maps  $L^2(0, T)$  into  $H_{(0)}^3 \times H_{(0)}^2$ .*

## 2.2 $C^1$ regularity of the end point map

Thanks to Proposition 2, we can consider the map  $\Theta_T$  defined by (10), and we know that it is continuous from  $L^2(0, T)$  to  $H_{(0)}^2 \times H^1(0, 1)$ . The goal of this section is the proof of the following hidden regularization effect.

**Theorem 3** *Let  $T > 0$  and  $\mu \in W^{2,\infty}(0, 1)$ . The map  $\Theta_T$  defined by (10) is  $C^1$  between the following spaces*

$$\Theta_T : L^2(0, T) \rightarrow H_{(0)}^3 \times H_{(0)}^2(0, 1).$$

Moreover, for every  $u, v \in L^2(0, T)$ , we have

$$d\Theta_T(u).v = \left(W, \frac{\partial W}{\partial t}\right)(T) \quad (22)$$

where  $W$  is the weak solution of

$$\begin{cases} \frac{\partial^2 W}{\partial t^2} = \frac{\partial^2 W}{\partial x^2} + u(t)\mu(x)W(t, x) + v(t)\mu(x)w(t, x), x \in (0, 1), t \in (0, T), \\ \frac{\partial W}{\partial x}(t, 0) = \frac{\partial W}{\partial x}(t, 1) = 0, \\ W(0, x) = 0, \\ \frac{\partial W}{\partial t}(0, x) = 0, \end{cases} \quad (23)$$

and  $w$  is the weak solution of (1), (6).

In Subsection 2.2.1, we state preliminary results useful for the proof of Theorem 3, which is detailed in Subsection 2.2.2.

### 2.2.1 Preliminaries

**Lemma 1** *Let  $T > 0$ . There exists  $C = C(T) > 0$  such that, for every  $g \in L^2(0, T)$ ,*

$$\left(\sum_{k \in \mathbb{N}} \left| \int_0^T g(t) e^{ik\pi t} dt \right|^2\right)^{1/2} \leq C \|g\|_{L^2(0, T)}.$$

**Proof of Lemma 1:** Let  $n \in \mathbb{N}^*$  be such that  $2(n-1) < T \leq 2n$ . Continuing  $g$  by zero on  $[T, 2n]$  and using the Bessel Parseval inequality, we get

$$\sum_{k \in \mathbb{N}} \left| \frac{1}{2n} \int_0^T g(t) e^{ik\pi t} dt \right|^2 \leq \frac{1}{2n} \int_0^T |g(t)|^2 dt.$$

Thus, Lemma 1 holds with  $C(T) := \sqrt{2n}$ .  $\square$

For  $s \geq 0$ , we use the spaces

$$h^s(\mathbb{N}^*, \mathbb{C}) := \left\{ a = (a_k)_{k \in \mathbb{N}^*} \in \mathbb{C}^{\mathbb{N}}; \sum_{k=1}^{\infty} |k^s a_k|^2 < +\infty \right\}$$

equipped with the norm

$$\|a\|_{h^s} := \left( \sum_{k=1}^{\infty} |k^s a_k|^2 \right)^{1/2}.$$

Thanks to Lemma 1, we have the following result.

**Lemma 2** Let  $T > 0$ . There exists  $C = C(T) > 0$  such that, for every  $w \in L^2(0, T)$ ,  $f \in C^0([0, T], H^2(0, 1))$ , the sequence  $S_0 = (S_{0,k})_{k \in \mathbb{N}^*}$  defined by

$$S_{0,k} := \int_0^T w(t) \langle f(t), \varphi_k \rangle e^{i\sqrt{\lambda_k}t} dt, \forall k \in \mathbb{N}^*,$$

belongs to  $h^2(\mathbb{N}^*, \mathbb{C})$  and

$$\|S_0\|_{h^2} \leq C \|w\|_{L^2} \|f\|_{C^0([0,T], H^2)}.$$

**Proof of Lemma 2:** Thanks to the equation  $A\varphi_k = \lambda_k \varphi_k$ , two integrations by part and the equalities  $\varphi_k(1) = (-1)^k \sqrt{2}$ ,  $\varphi_k(0) = \sqrt{2}$  (see (3)), we get the decomposition

$$\begin{aligned} S_{0,k} &= \frac{1}{\lambda_k} \int_0^T w(t) \langle Af(t), \varphi_k \rangle e^{i\sqrt{\lambda_k}t} dt \\ &\quad + \frac{(-1)^k \sqrt{2}}{\lambda_k} \int_0^T w(t) f'(t, 1) e^{i\sqrt{\lambda_k}t} dt \\ &\quad - \frac{\sqrt{2}}{\lambda_k} \int_0^T w(t) f'(t, 0) e^{i\sqrt{\lambda_k}t} dt, \end{aligned}$$

called  $S_0 = S_0^a + S_0^b + S_0^c$ . Thanks to (3) and Cauchy-Schwarz inequality, we have

$$\begin{aligned} \|S_0^a\|_{h^2} &= \left( \sum_{k=1}^{\infty} \left| k^2 \frac{1}{\lambda_k} \int_0^T w(t) \langle Af(t), \varphi_k \rangle e^{i\sqrt{\lambda_k}t} dt \right|^2 \right)^{1/2} \\ &\leq \frac{1}{\pi^2} \left( \sum_{k=1}^{\infty} \left( \int_0^T |w(t) \langle Af(t), \varphi_k \rangle| dt \right)^2 \right)^{1/2} \\ &\leq \frac{1}{\pi^2} \left( \sum_{k=1}^{\infty} \|w\|_{L^2}^2 \int_0^T |\langle Af(t), \varphi_k \rangle|^2 dt \right)^{1/2} \\ &\leq \frac{\sqrt{T}}{\pi^2} \|w\|_{L^2} \|Af\|_{C^0([0,T], L^2)} \\ &\leq \frac{\sqrt{T}}{\pi^2} \|w\|_{L^2} \|f\|_{C^0([0,T], H^2)}. \end{aligned}$$

Thanks to Lemma 1, there exists  $C = C(T) > 0$  such that

$$\begin{aligned} \|S_0^b\|_{h^2} &\leq C \|w(t) f'(t, 1)\|_{L^2} \leq C \|w\|_{L^2} \|f\|_{C^0([0,T], H^2)}, \\ \|S_0^c\|_{h^2} &\leq C \|w(t) f'(t, 0)\|_{L^2} \leq C \|w\|_{L^2} \|f\|_{C^0([0,T], H^2)}. \square \end{aligned}$$

### 2.2.2 Proof of Theorem 3

**Proof of Theorem 3:** Let  $T > 0$  and  $\mu \in W^{2,\infty}(0, 1)$ .

*First step:* We prove that  $\Theta_T$  indeed maps  $L^2(0, T)$  into  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$ .

Let  $u \in L^2(0, T)$  and  $w$  be the weak solution of (1), (6). Let

$$x_k := \langle w(T), \varphi_k \rangle + \frac{1}{i\sqrt{\lambda_k}} \left\langle \frac{\partial w}{\partial t}(T), \varphi_k \right\rangle, \forall k \in \mathbb{N}^*. \quad (24)$$

It is sufficient to prove that  $(x_k)_{k \in \mathbb{N}^*}$  belongs to  $h^3(\mathbb{N}^*, \mathbb{C})$ . From the formulation of a weak solution, we get

$$x_k = \frac{1}{i\sqrt{\lambda_k}} \int_0^T u(t) \langle \mu w(t), \varphi_k \rangle e^{i\sqrt{\lambda_k}(T-t)} dt, \forall k \in \mathbb{N}^*.$$

From Proposition 2, we know that

$$\left( w, \frac{\partial w}{\partial t} \right) \in C^0([0, T], H_{(0)}^2 \times H^1(0, 1)).$$

Thus  $\mu w \in C^0([0, T], H^2)$ , and Lemma 2 proves that  $(x_k)_{k \in \mathbb{N}^*}$  belongs to  $h^3(\mathbb{N}^*, \mathbb{C})$ .

*Second step:* We prove that the linear map  $v \mapsto W$  is continuous from  $L^2(0, T)$  to  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$ . Let  $u, v \in L^2(0, T)$  and  $w, W$  be the solutions of (1), (6) and (23). Let

$$X_k := \langle W(T), \varphi_k \rangle + \frac{1}{i\sqrt{\lambda_k}} \left\langle \frac{\partial W}{\partial t}(T), \varphi_k \right\rangle, \forall k \in \mathbb{N}^*.$$

It is sufficient to prove that  $X := (X_k)_{k \in \mathbb{N}^*}$  belongs to  $h^3(\mathbb{N}^*, \mathbb{C})$  and

$$\|X\|_{h^3} \leq C\|v\|_{L^2},$$

for some constant  $C = C(T, \mu, \|u\|_{L^2})$ . From the formulation of a weak solution, we get

$$X_k = \frac{1}{i\sqrt{\lambda_k}} \int_0^T \left( u(t) \langle \mu W(t), \varphi_k \rangle + v(t) \langle \mu w(t), \varphi_k \rangle \right) e^{i\sqrt{\lambda_k}(T-t)} dt, \forall k \in \mathbb{N}^*.$$

From Proposition 2, we know that

$$\left\| \left( W, \frac{\partial W}{\partial t} \right) \right\|_{C^0([0, T], H_{(0)}^2 \times H^1)} \leq C\|v\|_{L^2},$$

where  $C = C(T, \mu, \|v\|_{L^2})$ . Thus, applying Lemma 2, we get

$$\|X\|_{h^3} \leq C[\|u\|_{L^2} \|\mu W\|_{C^0([0, T], H^2)} + \|v\|_{L^2} \|\mu w\|_{C^0([0, T], H^2)}] \leq C\|v\|_{L^2}$$

where  $C = C(T, \mu, \|u\|_{L^2})$ .

*Third step:* We prove that  $\Theta_T : L^2(0, T) \rightarrow H_{(0)}^3 \times H_{(0)}^2(0, 1)$  is differentiable and (22) holds. Let  $u, v \in L^2(0, T)$ ,  $w, W, \tilde{w}$  be the weak solutions of (1), (6), (23) and

$$\begin{cases} \frac{\partial^2 \tilde{w}}{\partial t^2} = \frac{\partial^2 \tilde{w}}{\partial x^2} + (u + v)(t) \mu \tilde{w}, x \in (0, 1), t \in (0, T), \\ \frac{\partial \tilde{w}}{\partial x}(t, 0) = \frac{\partial \tilde{w}}{\partial x}(t, 1) = 0, \\ \tilde{w}(0, x) = 1, \\ \frac{\partial \tilde{w}}{\partial t}(0, x) = 0. \end{cases} \quad (25)$$

Then,  $\Delta := \tilde{w} - w - W$  is the weak solution of

$$\begin{cases} \frac{\partial^2 \Delta}{\partial t^2} = \frac{\partial^2 \Delta}{\partial x^2} + (u + v) \mu \Delta + v \mu W, \\ \frac{\partial \Delta}{\partial x}(t, 0) = \frac{\partial \Delta}{\partial x}(t, 1) = 0, \\ \Delta(0, x) = 0, \\ \frac{\partial \Delta}{\partial t}(0, x) = 0. \end{cases} \quad (26)$$

We want to prove that

$$\left\| \left( \Delta, \frac{\partial \Delta}{\partial t} \right) (T) \right\|_{H_{(0)}^3 \times H_{(0)}^2} = o(\|v\|_{L^2}) \text{ when } \|v\|_{L^2} \rightarrow 0.$$

Let

$$y_k := \langle \Delta(T), \varphi_k \rangle + \frac{1}{i\sqrt{\lambda_k}} \left\langle \frac{\partial \Delta}{\partial t}(T), \varphi_k \right\rangle, \forall k \in \mathbb{N}^*.$$

It is sufficient to prove that  $y := (y_k)_{k \in \mathbb{N}^*}$  satisfies  $\|y\|_{h^3} = O(\|v\|_{L^2}^2)$  when  $\|v\|_{L^2} \rightarrow 0$ . From the formulation of a weak solution, we get

$$y_k = \frac{1}{i\sqrt{\lambda_k}} \int_0^T \left( (u + v)(t) \langle \mu \Delta(t), \varphi_k \rangle + v(t) \langle \mu W(t), \varphi_k \rangle \right) e^{i\sqrt{\lambda_k}(T-t)} dt, \forall k \in \mathbb{N}^*.$$

From Proposition 2, we know that, when  $\|v\|_{L^2} \leq 1$ , we have

$$\begin{aligned} \left\| \left( W, \frac{\partial W}{\partial t} \right) \right\|_{C^0([0,T], H_{(0)}^2 \times H^1)} &\leq C \|v\mu w\|_{L^1((0,T), H^1)} \\ &\leq C \|v\|_{L^2} \|w\|_{C^0([0,T], H^1)} \\ &\leq C \|v\|_{L^2}, \\ \left\| \left( \Delta, \frac{\partial \Delta}{\partial t} \right) \right\|_{C^0([0,T], H_{(0)}^2 \times H^1)} &\leq C \|v\mu W\|_{L^1((0,T), H^1)} \\ &\leq C \|v\|_{L^2} \|W\|_{C^0([0,T], H^1)} \\ &\leq C \|v\|_{L^2}^2, \end{aligned}$$

where  $C = C(\mu, T, \|u\|_{L^2}) > 0$ . Thus, applying Lemma 2, we deduce that

$$\|y\|_{h^3} \leq \|u + v\|_{L^2} \|\mu \Delta\|_{C^0([0,T], H^2)} + \|v\|_{L^2} \|\mu W\|_{C^0([0,T], H^2)} \leq C \|v\|_{L^2}^2.$$

*Fourth step: We prove the continuity of the map*

$$\begin{aligned} d\Theta_T : L^2(0, T) &\rightarrow \mathcal{L}_c(L^2(0, T), H_{(0)}^3 \times H_{(0)}^2(0, 1)) \\ u &\mapsto d\Theta_T(u). \end{aligned}$$

Actually, we prove this map is locally Lipschitz. Let  $u, \tilde{u}, v \in L^2(0, T)$  with  $\|u - \tilde{u}\|_{L^2} < 1$  and  $w, W, \tilde{w}, \tilde{W}$  be the weak solutions of (1), (6), (23) and

$$\begin{cases} \frac{\partial^2 \tilde{w}}{\partial t^2} = \frac{\partial^2 \tilde{w}}{\partial x^2} + \tilde{u}\mu\tilde{w}, \\ \frac{\partial \tilde{w}}{\partial x}(t, 0) = \frac{\partial \tilde{w}}{\partial x}(t, 1) = 0, \\ \tilde{w}(0, x) = 1, \\ \frac{\partial \tilde{w}}{\partial t}(0, x) = 0, \end{cases} \quad \begin{cases} \frac{\partial^2 \tilde{W}}{\partial t^2} = \frac{\partial^2 \tilde{W}}{\partial x^2} + \tilde{u}\mu\tilde{W} + v\mu\tilde{w}, \\ \frac{\partial \tilde{W}}{\partial x}(t, 0) = \frac{\partial \tilde{W}}{\partial x}(t, 1) = 0, \\ \tilde{W}(0, x) = 0, \\ \frac{\partial \tilde{W}}{\partial t}(0, x) = 0. \end{cases}$$

We have

$$[d\Theta_T(u) - d\Theta_T(\tilde{u})].v = \left( \Xi, \frac{\partial \Xi}{\partial t} \right)(T)$$

where  $\Xi := W - \tilde{W}$  is the weak solution of

$$\begin{cases} \frac{\partial^2 \Xi}{\partial t^2} = \frac{\partial^2 \Xi}{\partial x^2} + u\mu\Xi + (u - \tilde{u})\mu\tilde{W} + v\mu(w - \tilde{w}), \\ \frac{\partial \Xi}{\partial x}(t, 0) = \frac{\partial \Xi}{\partial x}(t, 1) = 0, \\ \Xi(0, x) = 0, \\ \frac{\partial \Xi}{\partial t}(0, x) = 0. \end{cases}$$

Let

$$z_k := \langle \Xi(T), \varphi_k \rangle + \frac{1}{i\sqrt{\lambda_k}} \left\langle \frac{\partial \Xi}{\partial t}(T), \varphi_k \right\rangle, \forall k \in \mathbb{N}^*.$$

It is sufficient to prove that  $z := (z_k)_{k \in \mathbb{N}^*}$  satisfies

$$\|z\|_{h^3} \leq C \|u - \tilde{u}\|_{L^2} \|v\|_{L^2}, \quad (27)$$

where  $C = C(\mu, T, \|u\|_{L^2}) > 0$ . We have, for every  $k \in \mathbb{N}^*$ ,

$$z_k = \frac{1}{i\sqrt{\lambda_k}} \int_0^T \left( u(t) \langle \mu \Xi(t), \varphi_k \rangle + (u - \tilde{u})(t) \langle \mu \tilde{W}(t), \varphi_k \rangle + v(t) \langle \mu(w - \tilde{w})(t), \varphi_k \rangle \right) e^{i\sqrt{\lambda_k}(T-t)} dt.$$

Thus, applying Lemma 2, we get

$$\|z\|_{h^3} \leq C [\|u\|_{L^2} \|\Xi\|_{C^0([0,T], H^2)} + \|u - \tilde{u}\|_{L^2} \|\tilde{W}\|_{C^0([0,T], H^2)} + \|v\|_{L^2} \|w - \tilde{w}\|_{C^0([0,T], H^2)}],$$

where  $C = C(\mu, T, \|u\|_{L^2}) > 0$ . Thanks to Proposition 2, we have

$$\|w - \tilde{w}\|_{C^0([0,T], H_{(0)}^2)} \leq C \|(u - \tilde{u})\mu w\|_{L^1((0,T), H^1)} \leq C \|u - \tilde{u}\|_{L^2},$$

$$\begin{aligned}
\|\tilde{W}\|_{C^0([0,T],H_{(0)}^2)} &\leq C\|v\mu\tilde{w}\|_{L^1((0,T),H^1)} \leq C\|v\|_{L^2}, \\
\|\Xi\|_{C^0([0,T],H_{(0)}^2)} &\leq C\|(u-\tilde{u})\mu\tilde{W} + v\mu(w-\tilde{w})\|_{L^1((0,T),H^1)} \\
&\leq C[\|u-\tilde{u}\|_{L^2}\|\tilde{W}\|_{C^0([0,T],H^1)} + \|v\|_{L^2}\|w-\tilde{w}\|_{C^0([0,T],H^1)}] \\
&\leq C\|u-\tilde{u}\|_{L^2}\|v\|_{L^2},
\end{aligned}$$

where  $C = C(\mu, T, \|u\|_{L^2}) > 0$ . Therefore, we have (27).  $\square$

### 3 Controllability of the linearized system

The goal of this section is the proof of the following results.

**Theorem 4** *Let  $\mu \in W^{2,\infty}(0,1)$  be such that (5) holds.*

(1) *Let  $T > 2$ . The linear map  $d\Theta_T(0) : L^2(0,T) \rightarrow H_{(0)}^3 \times H_{(0)}^2(0,1)$  has a continuous right inverse  $d\Theta_T(0)^{-1} : H_{(0)}^3 \times H_{(0)}^2(0,1) \rightarrow L^2(0,T)$ .*

(2) *Let  $T = 2$ . The image of the linear map  $d\Theta_T(0) : L^2(0,T) \rightarrow H_{(0)}^3 \times H_{(0)}^2(0,1)$  is a vector subspace  $R_T$  of  $H_{(0)}^3 \times H_{(0)}^2(0,1)$  with codimension one, and there exists a continuous (left and right) inverse  $d\Theta_T(0)^{-1} : R_T \rightarrow L^2(0,T)$ .*

(3) *Let  $T < 2$ . The image of the linear map  $d\Theta_T(0) : L^2(0,T) \rightarrow H_{(0)}^3 \times H_{(0)}^2(0,1)$  is a vector subspace  $R_T$  of  $H_{(0)}^3 \times H_{(0)}^2(0,1)$  with infinite codimension, and there exists a continuous (left and right) inverse  $d\Theta_T(0)^{-1} : R_T \rightarrow L^2(0,T)$ .*

This section is organized as follows. In Subsection 3.1, we state preliminary results, useful for the proof of Theorem 4, which is detailed in Subsection 3.2.

#### 3.1 Preliminaries: trigonometric moment problems

Let us introduce the space

$$l_r^2([-1, +\infty), \mathbb{C}) := \{(d_k)_{k \geq -1}; d_{-1}, d_0 \in \mathbb{R}\}, \quad (28)$$

equipped with the norm

$$\|d\|_{l_r^2} := \left( \sum_{k=-1}^{\infty} |d_k|^2 \right)^{1/2}.$$

**Proposition 3** *Let  $T > 2$ . There exists a continuous linear map*

$$\begin{aligned}
L_T : l_r^2([-1, +\infty), \mathbb{C}) &\rightarrow L^2(0,T) \\
d = (d_k)_{k \geq -1} &\mapsto L_T(d)
\end{aligned}$$

*such that, for every sequence  $d = (d_k)_{k \geq -1} \in l_r^2([-1, +\infty), \mathbb{C})$  the function  $u := L_T(d)$  solves the moment problem*

$$\begin{cases} \int_0^T tu(t)dt = d_{-1}, \\ \int_0^T u(t)e^{ik\pi t}dt = d_k, \forall k \in \mathbb{N}. \end{cases} \quad (29)$$

**Proof of Proposition 3:** Let  $T > 2$ . The set

$$Z := \text{Cl}_{L^2((0,T),\mathbb{C})} \left( \text{Span}\{e^{ik\pi t}; k \in \mathbb{Z}\} \right)$$



(i.e.  $Z$  is the closure in  $L^2((0, T), \mathbb{C})$  of the vector space generated by the set  $\{e^{ik\pi t}; k \in \mathbb{Z}\}$ ) is a closed vector subspace of  $L^2((0, T), \mathbb{C})$  with infinite codimension. Let us prove that  $t \notin Z$ . Working by contradiction, we assume that  $t \in Z$ . With successive integrations, we get

$$t^j \in \text{Cl}_{C^0([0, T], \mathbb{C})}(\text{Span}\{t, e^{ik\pi t}; k \in \mathbb{Z}\}), \forall j \in \mathbb{N} \text{ with } j \geq 2.$$

The Stone Weierstrass theorem ensures that  $\{1, t^j; j \in \mathbb{N}, j \geq 2\}$  is dense in  $C^0([0, T], \mathbb{C})$ , thus, it is also dense in  $L^2((0, T), \mathbb{C})$ . Since  $t \in Z$ , we deduce that  $Z$  is dense in  $L^2((0, T), \mathbb{C})$ , which is impossible. Therefore,  $t \notin Z$ , and we have the following orthogonal decomposition

$$\begin{aligned} L^2((0, T), \mathbb{C}) &= Z \oplus Z^\perp \\ t &= z + z^\perp \end{aligned}$$

where  $z^\perp \neq 0$ . For  $d = (d_k)_{k \geq -1} \in l_r^2([-1, +\infty), \mathbb{C})$ , we define

$$L_T(d) := v + \left( d_{-1} - \int_0^T tv(t)dt \right) \frac{z^\perp}{\|z^\perp\|_{L^2}^2}$$

where

$$v := \left( \sum_{k \in \mathbb{Z}} d_k e^{-ik\pi t} \right) 1_{[0, 2]}(t)$$

and  $d_{-k} := \overline{d_k}, \forall k \in \mathbb{N}^*$ . The function  $L_T(u)$  is real valued because  $v$  and  $z^\perp$  are. From Bessel Parseval equation, we have

$$\|v\|_{L^2(0, T)}^2 = \frac{1}{2} \left[ |d_0|^2 + 2 \sum_{k=1}^{\infty} |d_k|^2 \right],$$

thus there exists  $C = C(T)$  such that

$$\|L_T(d)\|_{L^2(0, T)} \leq C(T) \|d\|_{l_r^2}. \square$$

The following proposition is a consequence of a more general result due to Horvath and Joo in [27].

**Proposition 4** *For every  $T \in (0, 2\pi)$ , there exists an extraction  $\xi : \mathbb{Z} \rightarrow \mathbb{Z}$  such that  $(e^{i\xi(k)t})_{k \in \mathbb{Z}}$  is a Riesz basis of  $L^2(0, T)$ .*

We also have the following stronger result, for particular values of  $T$ .

**Proposition 5** *Let  $T \in (0, 2\pi)$  of the form*

$$T = \frac{(2r-1)\pi}{p} \text{ with } r, p \in \mathbb{N}^*.$$

*There exists an extraction  $\xi : \mathbb{Z} \rightarrow \mathbb{Z}$  such that  $\xi(-k) = -\xi(k), \forall k \in \mathbb{Z}$  and  $(e^{i\xi(k)t})_{k \in \mathbb{Z}}$  is a Riesz basis of  $L^2(0, T)$ .*

**Proof of Proposition 5:** First, let us recall that the Kadec 1/4 Theorem says that, if the real valued sequence  $(\delta_n)_{n \in \mathbb{Z}}$  satisfies

$$\sup_{n \in \mathbb{Z}} |\delta_n| < 1/4,$$

then  $(e^{i(n+\delta_n)t})_{n \in \mathbb{Z}}$  is a Riesz basis of  $L^2(0, 2\pi)$ . Avdonin made the important remark that here, the 1/4 bound is sufficient to hold only for an average of the perturbations  $\delta_n$ . Namely, if

- $(\delta_n)_{n \in \mathbb{Z}}$  is bounded,
- $(n + \delta_n)_{n \in \mathbb{Z}}$  is separated, i.e.

$$\inf\{(n + \delta_n) - (m + \delta_m); n, m \in \mathbb{Z}, n \neq m\} > 0$$

- and we have

$$\lim_{K \rightarrow +\infty} \sup_{x \in \mathbb{R}} \frac{1}{K} \left| \sum_{x < n < x+K} \delta_n \right| < \frac{1}{4},$$

then  $(e^{i(n+\delta_n)t})_{n \in \mathbb{Z}}$  is a Riesz basis of  $L^2(0, 2\pi)$  (see [4]).

Now, let us prove Proposition 5. Let  $\xi : \mathbb{Z} \rightarrow \mathbb{Z}$  be the extraction such that  $\xi(0) = 0$  and the image of  $\xi$  is

$$R[\xi] = \cup_{n \in \mathbb{Z}} \{2np - r + 1, 2np - r + 2, \dots, 2np + r - 1\}.$$

This means that we keep  $(2r - 1)$  frequencies over  $2p$ , in chains centered at the frequencies  $2np$ ,  $n \in \mathbb{Z}$ . For this extraction, the average shift (with respect to  $\{2pn/(2r - 1); n \in \mathbb{Z}\}$ ) is equal to zero. Indeed, on any chain, the global shift is equal to zero. Thus,  $(e^{i\xi(k)t})_{k \in \mathbb{Z}}$  is a Riesz basis of  $L^2(0, T)$ .  $\square$

### 3.2 Study of the linearized system

The goal of this subsection is the proof of Theorem 4.

**Proof of Theorem 4:** Let  $\mu \in W^{2,\infty}(0, 1)$  be such that (5) holds. Let  $v \in L^2(0, T)$ . We have

$$d\Theta_T(0).v = \left( W, \frac{\partial W}{\partial t} \right) (T)$$

where  $W$  is the weak solution of

$$\begin{cases} \frac{\partial^2 W}{\partial t^2} = \frac{\partial^2 W}{\partial x^2} + v(t)\mu(x), x \in (0, 1), t \in (0, T), \\ \frac{\partial W}{\partial x}(t, 0) = \frac{\partial W}{\partial x}(t, 1) = 0, \\ W(0, x) = 0, \\ \frac{\partial W}{\partial t}(0, x) = 0. \end{cases} \quad (30)$$

We have

$$\begin{aligned} W(T) &= \left( \langle \mu, \varphi_0 \rangle \int_0^T (T - t)v(t)dt \right) \varphi_0 + \sum_{k=1}^{\infty} \left( \frac{\langle \mu, \varphi_k \rangle}{\sqrt{\lambda_k}} \int_0^T v(t) \sin[\sqrt{\lambda_k}(T - t)]dt \right) \varphi_k, \\ \frac{\partial W}{\partial t}(T) &= \left( \langle \mu, \varphi_0 \rangle \int_0^T v(t)dt \right) \varphi_0 + \sum_{k=1}^{\infty} \left( \langle \mu, \varphi_k \rangle \int_0^T v(t) \cos[\sqrt{\lambda_k}(T - t)]dt \right) \varphi_k. \end{aligned}$$

Thus, for  $(W_f, \dot{W}_f) \in H_{(0)}^3 \times H_{(0)}^2(0, 1)$ , the equality  $d\Theta_T(0).v = (W_f, \dot{W}_f)$  is equivalent to the moment problem

$$\begin{cases} \int_0^T (T - t)v(t)dt = d_{-1}(W_f, \dot{W}_f), \\ \int_0^T v(t)dt = d_0(W_f, \dot{W}_f), \\ \int_0^T v(t)e^{-i\sqrt{\lambda_k}t}dt = d_k(W_f, \dot{W}_f), \forall k \in \mathbb{N}^*, \end{cases}$$

where  $d(W_f, \dot{W}_f) = (d_k(W_f, \dot{W}_f))_{k \geq -1}$  is the sequence defined by

$$\begin{aligned} d_{-1}(W_k, \dot{W}_f) &:= \frac{\langle W_f, \varphi_0 \rangle}{\langle \mu, \varphi_0 \rangle}, \\ d_0(W_k, \dot{W}_f) &:= \frac{\langle \dot{W}_f, \varphi_0 \rangle}{\langle \mu, \varphi_0 \rangle}, \\ d_k(W_k, \dot{W}_f) &:= \frac{e^{-i\sqrt{\lambda_k}T}}{\langle \mu, \varphi_k \rangle} \left( \langle \dot{W}_f, \varphi_k \rangle + i\sqrt{\lambda_k} \langle W_f, \varphi_k \rangle \right), \forall k \in \mathbb{N}^*. \end{aligned} \quad (31)$$

Thanks to (5), the map

$$\begin{aligned} d : H_{(0)}^3 \times H_{(0)}^2(0, 1) &\rightarrow l_r^2([-1, +\infty), \mathbb{C}) \\ (W_f, \dot{W}_f) &\mapsto d(W_f, \dot{W}_f) \end{aligned}$$

is continuous (see (28) for a definition of  $l_r^2([-1, +\infty), \mathbb{C})$ ).

(1) We assume  $T > 2$ . Thanks to Proposition 3, the expression

$$d\Theta_T(0)^{-1}(W_f, \dot{W}_f) := L_T[d(W_f, \dot{W}_f)]$$

gives a suitable right inverse.

(2) We assume  $T = 2$ . Then the family  $(e^{ik\pi t})_{k \in \mathbb{Z}}$  is an orthonormal basis of  $L^2(0, T)$  and we have

$$(T - t) = \sum_{k \in \mathbb{Z}} \alpha_k e^{-ik\pi t} \text{ in } L^2(0, T),$$

where  $(\alpha_k)_{k \in \mathbb{Z}} \in l^2(\mathbb{Z}, \mathbb{C})$ . Then, the image of  $d\Theta_T(0)$  is the vector space

$$R_T := \left\{ (W_f, \dot{W}_f) \in H_{(0)}^3 \times H_{(0)}^2(0, 1); d_{-1}(W_f, \dot{W}_f) = \sum_{k \in \mathbb{Z}} \alpha_k \tilde{d}_k(W_f, \dot{W}_f) \right\},$$

where

$$\begin{aligned} \tilde{d}_k(W_f, \dot{W}_f) &:= d_k(W_f, \dot{W}_f), \forall k \in \mathbb{N}, \\ \tilde{d}_{-k}(W_f, \dot{W}_f) &:= d_k(W_f, \dot{W}_f), \forall k \in \mathbb{N}^*. \end{aligned} \tag{32}$$

The map  $d\Theta_T(0) : L^2(0, T) \rightarrow R_T$  has an inverse defined by

$$d\Theta_T(0)^{-1}(W_f, \dot{W}_f) = t \mapsto \sum_{k \in \mathbb{Z}} \tilde{d}_k(W_f, \dot{W}_f) e^{ik\pi t},$$

which is continuous from  $R_T$  (equipped with the  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$ -norm) to  $L^2(0, T)$ , thanks to the Bessel Parseval equality.

(3) We assume  $T < 2$ . Let  $\xi : \mathbb{Z} \rightarrow \mathbb{Z}$  be an extraction such that  $(e^{-i\xi(k)\pi t})_{k \in \mathbb{Z}}$  is a Riesz basis of  $L^2(0, T)$  (see Proposition 4). Then, there exists  $(\beta_k)_{k \in \mathbb{Z}} \in l^2(\mathbb{Z}, \mathbb{C})$  such that

$$T - t = \sum_{k \in \mathbb{Z}} \beta_k e^{-i\xi(k)\pi t} \text{ in } L^2(0, T)$$

and for every  $n \in \mathbb{N}$  that do not belong to the image of  $\xi$ , there exists  $(\gamma_k^n)_{k \in \mathbb{Z}} \in l^2(\mathbb{Z}, \mathbb{C})$  such that

$$e^{-in\pi t} = \sum_{k \in \mathbb{Z}} \gamma_k^n e^{-i\xi(k)\pi t} \text{ in } L^2(0, T).$$

Then, the image of  $d\Theta_T(0)$  is the vector space

$$\begin{aligned} R_T := \left\{ (W_f, \dot{W}_f) \in H_{(0)}^3 \times H_{(0)}^2(0, 1); d_{-1}(W_f, \dot{W}_f) = \sum_{k \in \mathbb{Z}} \beta_k \tilde{d}_{\xi(k)}(W_f, \dot{W}_f) \text{ and} \right. \\ \left. \forall n \in \mathbb{N} - R(\xi), d_n(W_f, \dot{W}_f) = \sum_{k \in \mathbb{Z}} \gamma_k^n \tilde{d}_{\xi(k)}(W_f, \dot{W}_f) \right\}. \end{aligned} \tag{33}$$

The set  $R_T$  is a vector subspace of  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$  with infinite codimension because it is defined by an infinite number of linearly independent relations. Let  $(\zeta_k)_{k \in \mathbb{Z}}$  be the biorthogonal family to  $(e^{-i\xi(k)\pi t})_{k \in \mathbb{Z}}$  in  $L^2(0, T)$ . Then, the map  $d\Theta_T(0) : L^2(0, T) \rightarrow R_T$  has a continuous inverse  $d\Theta_T(0)^{-1} : R_T \rightarrow L^2(0, T)$  defined by

$$d\Theta_T(0)^{-1}(W_f, \dot{W}_f) = \sum_{k \in \mathbb{Z}} \tilde{d}_{\xi(k)}(W_f, \dot{W}_f) \zeta_k. \square$$

## 4 Second order term

Using the same kind of arguments as in the proof of Theorem 3, one may prove the following result.

**Proposition 6** *Let  $\mu \in W^{2,\infty}(0,1)$  and  $T > 0$ . The map  $\Theta_T$  defined by (10) is twice differentiable at 0 and*

$$d^2\Theta_T(0).(v, v) = \left( \nu, \frac{\partial \nu}{\partial t} \right) (T)$$

where  $\nu$  is the weak solution of

$$\begin{cases} \frac{\partial \nu}{\partial t} = \frac{\partial^2 \nu}{\partial x^2} + v(t)\mu(x)W, \\ \frac{\partial \nu}{\partial x}(t, 0) = \frac{\partial \nu}{\partial x}(t, 1) = 0, \\ \nu(0, x) = 0, \\ \frac{\partial \nu}{\partial t}(0, x) = 0, \end{cases} \quad (34)$$

and  $W$  is the weak solution of (30).

The main result of this section is the following one.

**Proposition 7** *Let  $\mu \in W^{2,\infty}(0,1)$  be such that (5) and (7) hold and  $T \in (0, 2]$ . We assume that, either  $T = 2$ , or  $T = (2r - 1)/p$  with  $p, r \in \mathbb{N}^*$ . The image of the quadratic form  $d^2\Theta_T(0)$  is not contained in the image of the linear map  $d\Theta_T(0)$ .*

The following Lemma is useful for the proof of Proposition 7.

**Lemma 3** *Let  $T > 0$ ,  $D := \{(t, \tau) \in \mathbb{R}^2; 0 < \tau < t < T\}$  and  $h \in L^2(D, \mathbb{R})$ . If*

$$\int_0^T v(t) \int_0^t v(\tau) h(t, \tau) d\tau dt = 0, \forall v \in L^2(0, T),$$

then  $h = 0$ .

**Proof of Lemma 3:** We consider the quadratic form

$$\begin{aligned} Q : L^2(0, T) &\rightarrow \mathbb{R} \\ v &\mapsto Q(v) := \int_0^T v(t) \int_0^t v(\tau) h(t, \tau) d\tau dt. \end{aligned}$$

It is easy to prove that

$$\nabla Q(v) = t \mapsto \int_0^T v(\tau) \{h(t, \tau) 1_{\tau < t} + h(\tau, t) 1_{\tau > t}\} d\tau.$$

Since  $Q \equiv 0$ , we have  $\nabla Q \equiv 0$ , i.e.

$$\int_0^T v(\tau) \{h(t, \tau) 1_{\tau < t} + h(\tau, t) 1_{\tau > t}\} d\tau = 0, \text{ a.e. } t \in [0, T], \forall v \in L^2(0, T).$$

Thus,  $h(t, \tau) = 0$ , a.e.  $(t, \tau) \in D$ .  $\square$

**Proof of Proposition 7:** Let  $\mu \in W^{2,\infty}(0,1)$  be such that (5) and (7) hold and  $T \in (0, 2]$ . We assume that, either  $T = 2$ , or  $T = (2r - 1)/p$  with  $p, r \in \mathbb{N}^*$ .

*First step: Let us present the global strategy of the proof.* Let  $\xi : \mathbb{Z} \rightarrow \mathbb{Z}$  be such that

$$\xi(-k) = -\xi(k), \forall k \in \mathbb{N}^* \quad (35)$$

and  $(e^{-i\xi(k)\pi t})_{k \in \mathbb{Z}}$  is a Riesz basis of  $L^2(0, T)$  (see Proposition 5 for  $T < 2$  and take  $\xi(k) = k$  for  $T = 2$ ). There exists a unique sequence  $(\alpha_k)_{k \in \mathbb{Z}} \in l^2(\mathbb{Z}, \mathbb{C})$  such that

$$T - t = \Re \left[ \sum_{k=0}^{\infty} \alpha_k e^{-i\xi(k)\pi t} \right] \text{ in } L^2(0, T). \quad (36)$$

We have seen in the proof of Theorem 4 that the image  $R_T$  of the linear map  $d\Theta_T(0)$  is contained in the vector space

$$\tilde{R}_T := \left\{ (W_f, \dot{W}_f) \in H_{(0)}^3 \times H_{(0)}^2(0, 1); d_{-1}(W_f, \dot{W}_f) = \Re \left[ \sum_{k=0}^{\infty} \alpha_k d_{\xi(k)}(W_f, \dot{W}_f) \right] \right\},$$

where  $(d_k(W_f, \dot{W}_f))_{k \geq -1}$  is defined by (31). In order to prove Proposition 7, it is sufficient to prove that the image of the quadratic form  $d^2\Theta_T(0)$  is not contained in  $\tilde{R}_T$ .

*Second step: Let us state an equivalent property for “ $d^2\Theta_T(0).(v, v) \in \tilde{R}_T$ ”. Let  $v \in L^2(0, T)$  and  $W, \nu$  be the weak solutions of (30) and (34). We have*

$$W(t) = \left( \langle \mu, \varphi_0 \rangle \int_0^t (t - \tau) v(\tau) d\tau \right) \varphi_0 + \sum_{k=1}^{\infty} \left( \frac{\langle \mu, \varphi_k \rangle}{\sqrt{\lambda_k}} \int_0^t v(\tau) \sin[\sqrt{\lambda_k}(t - \tau)] d\tau \right) \varphi_k, \quad (37)$$

$$\begin{aligned} \nu(T) &= \left( \int_0^T (T - t) v(t) \langle \mu W(t), \varphi_0 \rangle dt \right) \varphi_0 \\ &+ \sum_{k=1}^{\infty} \left( \frac{1}{\sqrt{\lambda_k}} \int_0^T v(t) \langle \mu W(t), \varphi_k \rangle \sin[\sqrt{\lambda_k}(T - t)] dt \right) \varphi_k \end{aligned} \quad (38)$$

and

$$\begin{aligned} \frac{\partial \nu}{\partial t}(T) &= \left( \int_0^T v(t) \langle \mu W(t), \varphi_0 \rangle dt \right) \varphi_0 \\ &+ \sum_{k=1}^{\infty} \left( \int_0^T v(t) \langle \mu W(t), \varphi_k \rangle \cos[\sqrt{\lambda_k}(T - t)] dt \right) \varphi_k. \end{aligned} \quad (39)$$

Let us assume that  $d^2\Theta_T(0).(v, v) \in \tilde{R}_T$ . Then we have

$$\frac{\langle \nu(T), \varphi_0 \rangle}{\langle \mu, \varphi_0 \rangle} = \Re \left[ \sum_{k=0}^{\infty} \alpha_k \frac{e^{-i\xi(k)\pi T}}{\langle \mu, \varphi_{\xi(k)} \rangle} \left( \langle \dot{\nu}(T), \varphi_{\xi(k)} \rangle + i\sqrt{\lambda_{\xi(k)}} \langle \nu(T), \varphi_{\xi(k)} \rangle \right) \right]. \quad (40)$$

Thanks to (38) and (36), we have

$$\begin{aligned} \frac{\langle \nu(T), \varphi_0 \rangle}{\langle \mu, \varphi_0 \rangle} &= \frac{1}{\langle \mu, \varphi_0 \rangle} \int_0^T (T - t) v(t) \langle \mu W(t), \varphi_0 \rangle dt \\ &= \frac{1}{\langle \mu, \varphi_0 \rangle} \int_0^T \Re \left[ \sum_{k=0}^{\infty} \alpha_k e^{-i\xi(k)\pi t} \right] v(t) \langle \mu W(t), \varphi_0 \rangle dt \\ &= \int_0^T v(t) \Re \left[ \sum_{k=0}^{\infty} \alpha_k \frac{\langle \mu W(t), \varphi_0 \rangle}{\langle \mu, \varphi_0 \rangle} e^{-i\xi(k)\pi t} \right] dt. \end{aligned}$$

Thanks to (38) and (39), we have

$$\begin{aligned} &\frac{e^{-i\xi(k)\pi T}}{\langle \mu, \varphi_{\xi(k)} \rangle} \left( \langle \dot{\nu}(T), \varphi_{\xi(k)} \rangle + i\sqrt{\lambda_{\xi(k)}} \langle \nu(T), \varphi_{\xi(k)} \rangle \right) \\ &= \frac{e^{-i\xi(k)\pi T}}{\langle \mu, \varphi_{\xi(k)} \rangle} \int_0^T v(t) \langle \mu W(t), \varphi_{\xi(k)} \rangle e^{i\sqrt{\lambda_{\xi(k)}}(T-t)} dt \\ &= \int_0^T v(t) \frac{\langle \mu W(t), \varphi_{\xi(k)} \rangle}{\langle \mu, \varphi_{\xi(k)} \rangle} e^{-i\xi(k)\pi t} dt. \end{aligned}$$

Therefore, the equality (40) gives

$$\int_0^T v(t) \Re \left[ \sum_{k=0}^{\infty} \alpha_k \left( \frac{\langle \mu W(t), \varphi_0 \rangle}{\langle \mu, \varphi_0 \rangle} - \frac{\langle \mu W(t), \varphi_{\xi(k)} \rangle}{\langle \mu, \varphi_{\xi(k)} \rangle} \right) e^{-i\xi(k)\pi t} \right] dt = 0, \quad (41)$$

or, equivalently,

$$\int_0^T v(t) \Re \left[ \sum_{k=0}^{\infty} \alpha_k \left\langle \mu W(t), \frac{\varphi_0}{\langle \mu, \varphi_0 \rangle} - \frac{\varphi_{\xi(k)}}{\langle \mu, \varphi_{\xi(k)} \rangle} \right\rangle e^{-i\xi(k)\pi t} \right] dt = 0. \quad (42)$$

Noticing that

$$\left\langle \mu \varphi_0, \frac{\varphi_0}{\langle \mu, \varphi_0 \rangle} - \frac{\varphi_{\xi(k)}}{\langle \mu, \varphi_{\xi(k)} \rangle} \right\rangle = 0,$$

(because  $\varphi_0 = 1$ ) and using (37), we get

$$\begin{aligned} & \left\langle \mu W(t), \frac{\varphi_0}{\langle \mu, \varphi_0 \rangle} - \frac{\varphi_{\xi(k)}}{\langle \mu, \varphi_{\xi(k)} \rangle} \right\rangle \\ &= \sum_{j=1}^{\infty} \frac{\langle \mu, \varphi_j \rangle}{\sqrt{\lambda_j}} \int_0^t v(\tau) \sin[\sqrt{\lambda_j}(t - \tau)] d\tau \left\langle \mu \varphi_j, \frac{\varphi_0}{\langle \mu, \varphi_0 \rangle} - \frac{\varphi_{\xi(k)}}{\langle \mu, \varphi_{\xi(k)} \rangle} \right\rangle. \end{aligned}$$

Therefore, the equality (42) gives

$$\int_0^T v(t) \int_0^t v(\tau) h(t, t - \tau) d\tau dt = 0, \quad (43)$$

where, for every  $s \in \mathbb{R}$ ,  $t \in [0, T]$ ,

$$h(t, s) := \Re \left[ \sum_{k=0}^{\infty} \alpha_k e^{-i\xi(k)\pi t} \sum_{j=1}^{\infty} \frac{\langle \mu, \varphi_j \rangle}{\sqrt{\lambda_j}} \sin[\sqrt{\lambda_j}s] \left\langle \mu \varphi_j, \frac{\varphi_0}{\langle \mu, \varphi_0 \rangle} - \frac{\varphi_{\xi(k)}}{\langle \mu, \varphi_{\xi(k)} \rangle} \right\rangle \right].$$

*Third step: Let us prove that  $h \in C^0(\mathbb{R}_s, L^2(0, T)_t)$ .* We introduce the decomposition  $h = h_1 - h_2$  where

$$h_1(t, s) := \Re \left[ \sum_{k=0}^{\infty} \alpha_k e^{-i\xi(k)\pi t} \sum_{j=1}^{\infty} \frac{\langle \mu, \varphi_j \rangle}{\sqrt{\lambda_j}} \sin[\sqrt{\lambda_j}s] \frac{\langle \mu \varphi_j, \varphi_0 \rangle}{\langle \mu, \varphi_0 \rangle} \right],$$

$$h_2(t, s) := \Re \left[ \sum_{k=0}^{\infty} \alpha_k g_k(s) e^{-i\xi(k)\pi t} \right] \text{ and } g_k(s) := \sum_{j=1}^{\infty} \frac{\langle \mu, \varphi_j \rangle}{\sqrt{\lambda_j}} \sin[\sqrt{\lambda_j}s] \frac{\langle \mu \varphi_j, \varphi_{\xi(k)} \rangle}{\langle \mu, \varphi_{\xi(k)} \rangle}.$$

Using (36), we get

$$h_1(t, s) = \frac{(T - t)}{\langle \mu, \varphi_0 \rangle} f(s) \text{ where } f(s) := \sum_{j=1}^{\infty} \frac{\langle \mu, \varphi_j \rangle^2}{\sqrt{\lambda_j}} \sin[\sqrt{\lambda_j}s].$$

Explicit computations show that

$$\exists C > 0 \text{ such that } \frac{\langle \mu, \varphi_j \rangle^2}{\sqrt{\lambda_j}} \leq \frac{C}{j^5}, \forall j \in \mathbb{N}^*.$$

Thus  $f \in C^0(\mathbb{R}_s, \mathbb{R})$  and  $h_1 \in C^0(\mathbb{R}_s, L^2(0, T)_t)$ . Explicit computations show that

$$\exists C > 0 \text{ such that } \left| \frac{\langle \mu, \varphi_j \rangle \langle \mu \varphi_j, \varphi_K \rangle}{\sqrt{\lambda_j}} \right| \leq \frac{C}{j_*^3(j - K)_*^2}, \forall j, K \in \mathbb{N},$$

thus

$$|g_k(s)| \leq C \xi(k)^2 \sum_{j=1}^{\infty} \frac{1}{j^2(j - \xi(k))_*^2}, \forall k \in \mathbb{N}, \forall s \in \mathbb{R}.$$

The decomposition

$$\frac{1}{j^2(j-K)^2} = \frac{2}{K^3} \left( \frac{1}{j} - \frac{1}{j-K} \right) + \frac{1}{K^2} \left( \frac{1}{j^2} + \frac{1}{(j-K)^2} \right),$$

allows to prove that

$$\exists C > 0 \text{ such that } \sum_{j=1}^{\infty} \frac{1}{j^2(j-K)_*^2} \leq \frac{C}{K^2}, \forall K \in \mathbb{N},$$

thus,

$$\exists C > 0 \text{ such that } |g_k(s)| \leq C, \forall s \in \mathbb{R}, \forall k \in \mathbb{N}.$$

We have  $\alpha_k g_k(\sigma) \rightarrow \alpha_k g_k(s)$  when  $\sigma \rightarrow s$ , for every  $k \in \mathbb{N}$ . Moreover,  $|\alpha_k g_k(\sigma)| \leq C|\alpha_k|$ , for every  $k \in \mathbb{N}$ ,  $\sigma \in \mathbb{R}$ , and the sequence  $(|\alpha_k|)_{k \in \mathbb{N}}$  belongs to  $l^2(\mathbb{N})$ . The dominated convergence theorem ensures that the sequence  $(\alpha_k g_k(\sigma))_{k \in \mathbb{N}}$  converges to  $(\alpha_k g_k(s))_{k \in \mathbb{N}}$  in  $l^2(\mathbb{N}, \mathbb{C})$  when  $\sigma \rightarrow s$ . Since  $(e^{-i\xi(k)\pi t})_{k \in \mathbb{Z}}$  is a Riesz basis of  $L^2(0, T)$ , we deduce that  $h_2(\cdot, \sigma) \rightarrow h_2(\cdot, s)$  in  $L^2(0, T)$ , when  $\sigma \rightarrow s$ . This ends the proof of the third step.

*Fourth step: Now, let us prove Proposition 7.* Let us assume that  $d^2\Theta_T(0).(v, v) \in \tilde{R}_T$  for every  $v \in L^2(0, T)$ . Then, thanks to the second step, the equality (43) holds for every  $v \in L^2(0, T)$ . Moreover  $h \in L^2(D, \mathbb{R})$  (see the third step), so we can apply Lemma 3, which gives  $h = 0$  in  $L^2(D)$ . Therefore, we also have  $\partial h / \partial s = 0$  in the sense of distributions over  $D$ . In the sense of distributions on  $\mathbb{R}_s \times [0, T]_t$ , we have

$$\frac{\partial h}{\partial s}(t, s) = \Re \left[ \sum_{k=0}^{\infty} \alpha_k e^{-i\xi(k)\pi t} \sum_{j=1}^{\infty} \langle \mu, \varphi_j \rangle \cos[\sqrt{\lambda_j} s] \left\langle \mu \varphi_j, \frac{\varphi_0}{\langle \mu, \varphi_0 \rangle} - \frac{\varphi_{\xi(k)}}{\langle \mu, \varphi_{\xi(k)} \rangle} \right\rangle \right]. \quad (44)$$

Moreover, working as in the third step, one can prove that  $\frac{\partial h}{\partial s} \in C^0(\mathbb{R}_s, L^2(0, T)_t)$ . Thus, the equality (44) holds for every  $s \in \mathbb{R}$  in  $L^2(0, T)_t$ . In particular, with  $s = 0$ , we get

$$\Re \left[ \sum_{k=0}^{\infty} \alpha_k e^{-i\xi(k)\pi t} \sum_{j=1}^{\infty} \langle \mu, \varphi_j \rangle \left\langle \mu \varphi_j, \frac{\varphi_0}{\langle \mu, \varphi_0 \rangle} - \frac{\varphi_{\xi(k)}}{\langle \mu, \varphi_{\xi(k)} \rangle} \right\rangle \right] = 0 \text{ in } L^2(0, T)_t,$$

or, equivalently,

$$\Re \left[ \sum_{k=0}^{\infty} \alpha_k e^{-i\xi(k)\pi t} \left\langle \mu^2, \frac{\varphi_0}{\langle \mu, \varphi_0 \rangle} - \frac{\varphi_{\xi(k)}}{\langle \mu, \varphi_{\xi(k)} \rangle} \right\rangle \right] = 0, \text{ in } L^2(0, T)_t.$$

But  $(e^{-i\xi(k)\pi t})_{k \in \mathbb{Z}}$  is a Riesz basis of  $L^2(0, T)$  and (35) holds, thus the previous equality implies

$$\alpha_k \left\langle \mu^2, \frac{\varphi_0}{\langle \mu, \varphi_0 \rangle} - \frac{\varphi_{\xi(k)}}{\langle \mu, \varphi_{\xi(k)} \rangle} \right\rangle = 0, \forall k \in \mathbb{N}. \quad (45)$$

Let us assume temporarily that the number of integers  $p \in \mathbb{N}$  such that

$$\left\langle \mu^2, \frac{\varphi_0}{\langle \mu, \varphi_0 \rangle} - \frac{\varphi_p}{\langle \mu, \varphi_p \rangle} \right\rangle = 0 \quad (46)$$

is finite. Then the equality (45) implies that only a finite number of  $\alpha_k$  may be different from zero. But this is in contradiction with (36) because, for every  $N \in \mathbb{N}$ , the family  $\{t, e^{ik\pi t}; -N \leq k \leq N\}$  is linearly independent. Therefore, the image of  $d^2\Theta_T(0)$  is not contained in  $\tilde{R}_T$ .

Now let us prove that the number of integers  $p \in \mathbb{N}$  such that (46) holds is finite. Integrations by parts give

$$\begin{aligned}\langle \mu, \varphi_p \rangle &= \frac{\sqrt{2}}{(p\pi)^2} \left( (-1)^p \mu'(1) - \mu'(0) \right) - \frac{\sqrt{2}}{(p\pi)^2} \int_0^1 \mu''(x) \cos(p\pi x) dx, \\ \langle \mu^2, \varphi_p \rangle &= \frac{\sqrt{2}}{(p\pi)^2} \left( (-1)^p (\mu^2)'(1) - (\mu^2)'(0) \right) - \frac{\sqrt{2}}{(p\pi)^2} \int_0^1 (\mu^2)''(x) \cos(p\pi x) dx.\end{aligned}$$

Since  $\mu'(1) \pm \mu'(0) \neq 0$ , we have

$$\frac{\langle \mu^2, \varphi_p \rangle}{\langle \mu, \varphi_p \rangle} \sim \frac{(-1)^p (\mu^2)' - (\mu^2)'(0)}{(-1)^p \mu'(1) - \mu'(0)} \text{ when } p \rightarrow +\infty.$$

Thus the assumption (7) implies that the number of integers  $p \in \mathbb{N}$  such that (46) holds is finite.  $\square$

## 5 Proof of Theorem 1

**Proof of Theorem 1:** Let  $\mu \in W^{2,\infty}(0,1)$  be such that (5) holds.

(1) Let  $T > 2$ . The map

$$\Theta_T : L^2(0, T) \rightarrow H_{(0)}^3 \times H_{(0)}^2(0, 1)$$

is  $C^1$  (see Theorem 3), and  $d\Theta_T(0)$  has a continuous right inverse (see Theorem 4 (1))

$$d\Theta_T(0)^{-1} : H_{(0)}^3 \times H_{(0)}^2(0, 1) \rightarrow L^2(0, T).$$

Thus, thanks to the inverse mapping theorem,  $\Theta_T$  has a local  $C^1$  right inverse.

(3) Let  $T < 2$ . First, let us assume that  $T = (2r - 1)/p$  with  $p, r \in \mathbb{N}^*$ . The set  $R_T$  defined by (33) is a closed vector subspace of the Hilbert space  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$ . Thus, we have the orthogonal decomposition

$$H_{(0)}^3 \times H_{(0)}^2(0, 1) = R_T \oplus R_T^\perp.$$

We consider the map

$$\begin{aligned}F_T : L^2(0, T) \times R_T^\perp &\rightarrow H_{(0)}^3 \times H_{(0)}^2(0, 1) \\ (u, y) &\mapsto \Theta_T(u) + y.\end{aligned}$$

Thanks to Theorem 3,  $F_T$  is  $C^1$ . Thanks to Theorem 4 (3), the continuous linear map

$$dF_T(0, 0) : L^2(0, T) \times R_T^\perp \rightarrow H_{(0)}^3 \times H_{(0)}^2(0, 1)$$

has a continuous inverse

$$dF_T(0, 0)^{-1} : H_{(0)}^3 \times H_{(0)}^2(0, 1) \rightarrow L^2(0, T) \times R_T^\perp$$

defined by

$$dF_T(0, 0)^{-1} \cdot (W_f, \dot{W}_f) := (d\Theta_T(0)^{-1} \cdot \mathcal{P}_{R_T}(W_f, \dot{W}_f), \mathcal{P}_{R_T^\perp}(W_f, \dot{W}_f))$$

where  $\mathcal{P}_{R_T}$  (resp.  $\mathcal{P}_{R_T^\perp}$ ) is the orthogonal projection from  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$  to  $R_T$  (resp.  $R_T^\perp$ ). Thanks to the inverse mapping theorem, the map  $F_T$  has a local inverse: there exists  $\delta, r > 0$  and a  $C^1$  map

$$F_T^{-1} : \mathcal{V}_T \rightarrow L^2(0, T) \times R_T^\perp$$



where

$$\mathcal{V}_T := \{(w_f, \dot{w}_f) \in H_{(0)}^3 \times H_{(0)}^2(0, 1); \|w_f - 1\|_{H_{(0)}^3} + \|\dot{w}_f\|_{H_{(0)}^2} < \delta\}$$

such that  $F_T^{-1}(1, 0) = (0, 0)$ ,  $F_T^{-1}[F_T(u, y)] = (u, y)$ , for every  $(u, y) \in L^2(0, T) \times R_T^\perp$  with  $\|u\|_{L^2} + \|y\|_{H_{(0)}^3 \times H_{(0)}^2} < r$  and  $F_T[F_T^{-1}(z)] = z$ , for every  $z \in \mathcal{V}_T$ . Let us denote by  $G_T$  the second component of  $F_T^{-1}$ . Then, the map

$$G_T : \mathcal{V}_T \rightarrow R_T^\perp$$

is locally surjective and we have

$$G_T[\Theta_T(u)] = 0, \forall u \in L^2(0, T) \text{ with } \|u\|_{L^2} < r.$$

This proves that the image of  $\Theta_T$  is locally a  $C^1$ -submanifold of  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$  with infinite codimension. This submanifold does not coincide with its tangent space at  $(1, 0)$  thanks to Proposition 7.

Now, let us consider an arbitrary  $T \in (0, 2)$ . Let  $T' \in (T, 2)$  be such that  $T' = (2r - 1)/p$  for some  $p, r \in \mathbb{N}^*$ . We continue the controls (defined on  $(0, T)$ ) by zero on  $(T, T')$ . Applying the previous result, we get

$$G_{T'}[e^{A(T'-T)}\Theta_T(u)] = 0, \forall u \in B_r[L^2(0, T)].$$

Thus, the map  $G_T := G_{T'} \circ e^{A(T'-T)}$  gives the conclusion.

**(2)** Let  $T = 2$ . First, let us prove that the nonlinear system is locally controllable up to codimension one. We consider the map

$$\begin{aligned} \tilde{\Theta}_T : L^2(0, T) &\rightarrow \tilde{\mathcal{V}}_T \\ u &\mapsto \left( w(T) - \int_0^1 w(T, x) dx, \frac{\partial w}{\partial t}(T) \right), \end{aligned}$$

where

$$\tilde{\mathcal{V}}_T := \left\{ (\tilde{w}_f, \dot{\tilde{w}}_f) \in H_{(0)}^3 \times H_{(0)}^2(0, 1); \int_0^1 \tilde{w}_f(x) dx = 0 \right\}.$$

Thanks to Theorem 3,  $\tilde{\Theta}_T$  is  $C^1$ . Thanks to Theorem 4 **(2)**, the continuous linear map

$$d\tilde{\Theta}_T(0) : L^2(0, T) \rightarrow \tilde{\mathcal{V}}_T$$

has a continuous inverse

$$d\tilde{\Theta}_T(0)^{-1} : \tilde{\mathcal{V}}_T \rightarrow L^2(0, T).$$

Thanks to the inverse mapping theorem,  $\tilde{\Theta}_T$  has a local  $C^1$  inverse. This proves the local controllability up to codimension one of (1) in time  $T = 2$ , in  $H_{(0)}^3 \times H_{(0)}^2(0, 1)$ , with  $L^2(0, T)$ -controls.

Working as in the proof of **(3)**, we get a locally surjective  $C^1$  map

$$G_T : H_{(0)}^3 \times H_{(0)}^2 \rightarrow \mathbb{R}$$

such that, for every  $u \in L^2(0, T)$  small enough,  $G_T[\Theta_T(u)] = 0$ . Thus, the image of  $\Theta_T$  is a  $C^1$ -submanifold of  $H_{(0)}^3 \times H_{(0)}^2$  with codimension one. Thanks to Proposition 7, this submanifold does not coincide with its tangent space at  $(1, 0)$ .  $\square$

## 6 Conclusion, open problems, perspectives

### 6.1 Same system, other reference trajectory

In this article, we have studied the local controllability of the system (1) around the reference trajectory

$$(w_{ref}(t, x) = 1, u_{ref}(t) = 0). \quad (47)$$

One may study the local controllability of the same system around other reference trajectories, for example

$$(w_{ref}(t, x) := \sin(K\pi t)\varphi_K(x), u_{ref}(t) = 0) \text{ for } K \in \mathbb{N}^*. \quad (48)$$

Let us explain why this problem is more difficult than the one solved in this article. The difficulty relies in the controllability of the linearized system. The linearized system of (1) around the reference trajectory (48) is

$$\begin{cases} \frac{\partial^2 W}{\partial t^2} = \frac{\partial^2 W}{\partial x^2} + u(t)\mu(x)w_{ref}(t, x), x \in (0, 1), t \in (0, T), \\ \frac{\partial W}{\partial x}(t, 0) = \frac{\partial W}{\partial x}(t, 1) = 0, \\ W(0, x) = \frac{\partial W}{\partial t}(0, x) = 0. \end{cases} \quad (49)$$

Working as in the proof of Theorem 4, one may prove that, for every  $(W_f, \dot{W}_f) \in H_{(0)}^3 \times H_{(0)}^2(0, 1)$ , the equality

$$\left(W, \frac{\partial W}{\partial t}\right)(T) = (W_f, \dot{W}_f)$$

is equivalent to the moment problem

$$\begin{aligned} \int_0^T (T-t)v(t)\sin(K\pi t)dt &= d_{-1}(W_f, \dot{W}_f), \\ \int_0^T v(t)\sin(K\pi t)e^{-ik\pi t}dt &= d_k(W_f, \dot{W}_f), \forall k \in \mathbb{N}, \end{aligned} \quad (50)$$

where  $(d_k(W_f, \dot{W}_f))_{k \geq -1}$  is defined by

$$\begin{aligned} d_{-1}(W_f, \dot{W}_f) &:= \frac{\langle W_f, \varphi_0 \rangle}{\langle \mu \varphi_K, \varphi_0 \rangle}, \\ d_k(W_f, \dot{W}_f) &:= \frac{e^{-i\sqrt{\lambda_k}T}}{\langle \mu \varphi_K, \varphi_k \rangle} \left( \langle \dot{W}_f, \varphi_k \rangle + i\sqrt{\lambda_k} \langle W_f, \varphi_k \rangle \right), \forall k \in \mathbb{N}. \end{aligned}$$

In order to apply the same strategy as in this article, one would need to prove that the family  $\{t \sin(K\pi t), \sin(K\pi t)e^{-ik\pi t}; k \in \mathbb{Z}\}$  satisfies the Riesz-basis property in  $L^2(0, T)$ , for every  $T > 2$ . However, this is false. Thus, the study of the local controllability of (1) around the reference trajectory (48) needs additional tools.

The same problem appears with Dirichlet boundary conditions, instead of Neumann boundary conditions in (1).

### 6.2 Other bilinear wave equations

Let us consider the following generalizations of the system (1).

$$\begin{cases} \frac{\partial^2 w}{\partial t^2}(t, x) = \frac{\partial^2 w}{\partial x^2}(t, x) + u(t) \left( \mu_1(x)w(t, x) + \mu_2(x) \frac{\partial w}{\partial x}(t, x) \right), x \in (0, 1), t \in (0, T), \\ \frac{\partial w}{\partial x}(t, 0) = \frac{\partial w}{\partial x}(t, 1) = 0, \end{cases} \quad (51)$$

$$\begin{cases} \frac{\partial^2 w}{\partial t^2}(t, x) = \frac{\partial^2 w}{\partial x^2}(t, x) + u(t) \left( \mu_1(x)w(t, x) + \mu_2(x) \frac{\partial w}{\partial t}(t, x) \right), x \in (0, 1), t \in (0, T), \\ \frac{\partial w}{\partial x}(t, 0) = \frac{\partial w}{\partial x}(t, 1) = 0. \end{cases} \quad (52)$$

Let us study the local controllability of (51) and (52) around the reference trajectory (47).

The operators

$$D(\mathcal{B}_1) := L^2 \times L^2(0, 1), \quad \mathcal{B}_1 \begin{pmatrix} 0 & 0 \\ \mu_1 \text{Id} + \mu_2 \partial_x & 0 \end{pmatrix}$$

$$D(\mathcal{B}_2) := L^2 \times L^2(0, 1), \quad \mathcal{B}_2 \begin{pmatrix} 0 & 0 \\ \mu_1 & \mu_2 \end{pmatrix}$$

are bounded on  $H_{(0)}^2 \times H^1(0, 1)$  when  $\mu_1, \mu_2 \in W^{2,\infty}(0, 1)$ . Thus Ball, Marsden and Slemrod's Theorem 2 applies: the systems (51) and (52) are not controllable in  $H_{(0)}^2 \times H^1(0, 1)$  with controls in  $L_{loc}^r[0, +\infty)$ ,  $r > 1$ . The Proposition 2 also holds with (1) replaced by (51) or (52): the Cauchy problems are well posed in  $H_{(0)}^2 \times H^1(0, 1)$  with controls in  $L_{loc}^1[0, +\infty)$ . The linearized system of (51) or (52) around the reference trajectory (47) is exactly the linearized system of (1) around the same trajectory. Thus, Theorem 4 may be used.

A new difficulty arises when one tries to prove the  $C^1$  regularity of the end point map (i.e. the adaptation of Theorem 3). Let us explain why. As in the proof of Theorem 3, we consider  $v \in L^2(0, T)$  and the weak solution  $w$  of (51) (resp. (52)) and (6). We want to prove that

$$\left(w, \frac{\partial w}{\partial t}\right)(T) \in H_{(0)}^3 \times H_{(0)}^2(0, 1).$$

Let  $x := (x_k)_{k \in \mathbb{N}^*}$  be defined by (24) and let us try to prove that this sequence belongs to  $h^3(\mathbb{N}^*, \mathbb{C})$ . From the formulation of a weak solution, we get

$$x_k = \frac{1}{i\sqrt{\lambda_k}} \int_0^T u(t) \left\langle \left( \mu_1 w + \mu_2 \frac{\partial w}{\partial x} \right)(t), \varphi_k \right\rangle e^{i\sqrt{\lambda_k}(T-t)} dt, \forall k \in \mathbb{N}^*$$

(resp.  $x_k = \frac{1}{i\sqrt{\lambda_k}} \int_0^T u(t) \left\langle \left( \mu_1 w + \mu_2 \frac{\partial w}{\partial t} \right)(t), \varphi_k \right\rangle e^{i\sqrt{\lambda_k}(T-t)} dt, \forall k \in \mathbb{N}^*$ ).

From Proposition 2, we know that

$$\left(w, \frac{\partial w}{\partial t}\right) \in C^0([0, T], H_{(0)}^2 \times H^1(0, 1)).$$

Thus  $\mu w \in C^0([0, T], H^2)$  and  $\mu \frac{\partial w}{\partial x} \in C^0([0, T], H^1)$  (resp.  $\mu \frac{\partial w}{\partial t} \in C^0([0, T], H^1)$ ). This regularity is not sufficient to apply Lemma 2. Thus, for the systems (51) and (52), if the end point map  $\Theta_T$  is  $C^1$  from  $L^2(0, T)$  to  $H_{(0)}^3 \times H_{(0)}^2(0, T)$ , then the proof of this property would be different from the one of Theorem 3.

Therefore, the local controllability of systems (51) and (52) around the reference trajectory (47) is an open problem.

### 6.3 Conjecture for 2D and 3D bilinear Schrödinger equations.

As emphasized in Subsection 1.3, the system (1) is a toy model for 2D Schrödinger equations, with bilinear controls (16). For such systems, we have the Weyl formula (17). We conjecture that, under generic assumptions on  $\mu$ ,

- for every  $T > 2\pi/d$ , the system (16) is locally exactly controllable around the ground state (or any eigenstate) in some function space (to be defined),
- for every  $T < 2\pi/d$ , the system (16) is not locally exactly controllable around the ground state: the reachable set is contained in a non flat submanifold of some functional space (to be defined), with infinite codimension.

Similarly, for 3D Schrödinger equations with bilinear control (i.e. equation (16) with  $\Omega$  a bounded open subset of  $\mathbb{R}^3$ ), we conjecture that, for every  $T > 0$ , the reachable set is a non flat submanifold of some functional space, with infinite codimension.

## A Genericity of the assumption on $\mu$

The goal of this section is the proof of the following result.

**Proposition 8** *The set  $\{\mu \in W^{2,\infty}(0,1); (5) \text{ and } (7) \text{ hold}\}$  is dense in  $W^{2,\infty}(0,1)$ .*

The following lemma will be useful in the proof of Proposition 8.

**Lemma 4** *Let  $\Phi_{\pm}, \Phi : W^{2,\infty}(0,1) \rightarrow \mathbb{R}$  be defined by  $\Phi(\mu) := \Phi_+(\mu)\Phi_-(\mu)$  and*

$$\Phi_{\pm}(\mu) := [(\mu^2)'(1) \pm (\mu^2)'(0)] \int_0^1 \mu(x) dx - [\mu'(1) \pm \mu'(0)] \int_0^1 \mu(x)^2 dx.$$

*For every  $\mu \in W^{2,\infty}(0,1)$  such that  $\mu'(1) \pm \mu'(0) \neq 0$  and  $\Phi(\mu) = 0$ , we have either  $d\Phi(\mu) \neq 0$  or  $d\Phi(\mu) = 0$  and  $d^2\Phi(\mu) \neq 0$ .*

**Proof of Lemma 4:** For every  $\mu, \nu \in W^{2,\infty}(0,1)$ , we have

$$\begin{aligned} d\Phi_{\pm}(\mu) \cdot \nu &= 2[(\mu\nu)'(1) \pm (\mu\nu)'(0)] \int_0^1 \mu + [(\mu^2)'(1) \pm (\mu^2)'(0)] \int_0^1 \nu \\ &\quad - [\nu'(1) \pm \nu'(0)] \int_0^1 \mu^2 - [\mu'(1) \pm \mu'(0)] \int_0^1 2\mu\nu. \end{aligned}$$

In particular, for every  $\nu \in C_c^\infty(0,1)$  such that  $\int_0^1 \nu = 0$ , we have

$$d\Phi_{\pm}(\mu) \cdot \nu = -2[\mu'(1) \pm \mu'(0)] \int_0^1 \mu\nu.$$

Let  $\mu \in W^{2,\infty}(0,1)$  be such that  $\mu'(1) \pm \mu'(0) \neq 0$  and  $\Phi(\mu) = 0$ .

*First case:* We assume  $\Phi_+(\mu) = 0$  and  $\Phi_-(\mu) \neq 0$ . Then, for every  $\nu \in C_c^\infty(0,1)$  such that  $\int_0^1 \nu = 0$  and  $\int_0^1 \mu\nu \neq 0$ , we have

$$d\Phi(\mu) \cdot \nu = [d\Phi_+(\mu) \cdot \nu] \Phi_-(\mu) = -2\Phi_-(\mu) [\mu'(1) + \mu'(0)] \int_0^1 \mu\nu \neq 0.$$

The case  $\Phi_-(\mu) = 0$  and  $\Phi_+(\mu) \neq 0$  may be treated similarly.

*Second case:* We assume  $\Phi_+(\mu) = \Phi_-(\mu) = 0$ . Then,  $d\Phi(\mu) = 0$  and, for every  $\nu \in C_c^\infty(0,1)$  such that  $\int_0^1 \nu = 0$  and  $\int_0^1 \mu\nu \neq 0$ , we have

$$\begin{aligned} d^2\Phi(\mu) \cdot \nu &= [d\Phi_+(\mu) \cdot \nu][d\Phi_-(\mu) \cdot \nu] \\ &= 4[\mu'(1) - \mu'(0)][\mu'(1) + \mu'(0)] \left( \int_0^1 \mu\nu \right)^2 \neq 0. \square \end{aligned}$$

**Proof of Proposition 8:** First, let us notice that

$$\mathcal{W} := \{\mu \in W^{2,\infty}(0,1); \mu'(0) \pm \mu'(1) \neq 0\}$$

is a dense open subset of  $W^{2,\infty}(0,1)$ . Thanks to Lemma 4, the set

$$\mathcal{V} := \{\mu \in \mathcal{W}; (7) \text{ holds}\}$$

is a dense open subset of  $\mathcal{W}$ . Now, let us prove that the set

$$\mathcal{U} := \{\mu \in \mathcal{V}; \langle \mu, \varphi_k \rangle \neq 0, \forall k \in \mathbb{N}\}$$

is dense in  $\mathcal{V}$ . For  $n \in \mathbb{N}$ , we introduce the set

$$\mathcal{U}_n := \{\mu \in \mathcal{V}; \langle \mu, \varphi_k \rangle \neq 0, \forall k \in \{0, \dots, n\}\},$$

with the convention  $\mathcal{U}_{-1} = \mathcal{V}$ . Then, the sequence  $\mathcal{U}_n$  is decreasing and

$$\mathcal{U} = \bigcap_{n=-1}^{\infty} \mathcal{U}_n.$$

We apply Baire Lemma : it is sufficient to prove that, for every  $n \geq -1$ ,  $\mathcal{U}_{n+1}$  is dense in  $\mathcal{U}_n$  for the  $W^{2,\infty}(0,1)$ -topology. Let  $n \geq -1$  and let  $\mu \in \mathcal{U}_n - \mathcal{U}_{n+1}$ . Then  $\langle \mu \varphi_1, \varphi_k \rangle \neq 0$  for  $k = 0, \dots, n$  and  $\langle \mu \varphi_1, \varphi_{n+1} \rangle = 0$ . There exists  $\epsilon^* > 0$  such that, for every  $\epsilon \in (0, \epsilon^*)$ ,  $\mu + \epsilon x^2 \in \mathcal{V}$ , because  $\mathcal{V}$  is an open subset of  $W^{2,\infty}(0,1)$ . Thanks to (9),  $\mu + \epsilon x^2 \in \mathcal{U}_{n+1}$  for every  $\epsilon \in (0, \epsilon^*)$  such that

$$\epsilon \neq -\frac{\langle \mu \varphi_1, \varphi_j \rangle}{\langle x^2 \varphi_1, \varphi_j \rangle}, \forall j \in \{0, \dots, n\}.$$

Thus  $\mathcal{U}_{n+1}$  is dense in  $\mathcal{U}_n$ . We have proved that  $\mathcal{U}$  is dense in  $W^{2,\infty}(0,1)$ .

Now, let us emphasize that

$$\mathcal{U} \subset \{\mu \in W^{2,\infty}(0,1); (5) \text{ and } (7) \text{ hold}\}.$$

Indeed, for  $\mu \in \mathcal{U}$ , and  $k \in \mathbb{N}^*$ , integrations by parts give (8). Since  $\mu'(0) \pm \mu'(1) \neq 0$ , there exists  $N \in \mathbb{N}$  such that, for every  $k \geq N$ ,

$$|\langle \mu, \varphi_k \rangle| \geq \frac{1}{(k\pi)^2} \max\{|\mu'(1) + \mu'(0)|, |\mu'(1) - \mu'(0)|\}.$$

Since  $\langle \mu, \varphi_k \rangle \neq 0, \forall k \in \mathbb{N}$ , there exists  $c > 0$  such that

$$|\langle \mu, \varphi_k \rangle| \geq \frac{c}{k_*^2}, \forall k \in \mathbb{N}. \square$$

**Acknowledgments:** The author thanks Jean-Michel Coron and Alain Haraux for interesting discussions and Miklos Horvath for helpful remarks, for the proof of Proposition 4 and the references [4] and [27].

## Contents

<b>1</b>	<b>Introduction</b>	<b>1</b>
1.1	Main result	1
1.2	Sketch of the proof	4
1.3	A brief bibliography	5
1.3.1	A previous negative result for this equation	5
1.3.2	Iterated Lie brackets for general bilinear systems	7
1.3.3	Wave equation with bilinear control	8
1.3.4	Wave equation with linear controls	8
1.3.5	Other results about infinite dimensional bilinear systems	8
1.4	A toy model for 2D quantum systems	9
1.5	Structure of this article	10
<b>2</b>	<b>Well posedness and <math>C^1</math> regularity of the end point map</b>	<b>10</b>
2.1	Existence, uniqueness, regularity and bounds	10
2.2	$C^1$ regularity of the end point map	12
2.2.1	Preliminaries	12
2.2.2	Proof of Theorem 3	13

<b>3</b>	<b>Controllability of the linearized system</b>	<b>16</b>
3.1	Preliminaries: trigonometric moment problems . . . . .	16
3.2	Study of the linearized system . . . . .	18
<b>4</b>	<b>Second order term</b>	<b>20</b>
<b>5</b>	<b>Proof of Theorem 1</b>	<b>24</b>
<b>6</b>	<b>Conclusion, open problems, perspectives</b>	<b>26</b>
6.1	Same system, other reference trajectory . . . . .	26
6.2	Other bilinear wave equations . . . . .	26
6.3	Conjecture for 2D and 3D bilinear Schrödinger equations. . . . .	27
<b>A</b>	<b>Genericity of the assumption on <math>\mu</math></b>	<b>28</b>

## References

- [1] R. Adami and U. Boscain. Controllability of the Schroedinger Equation via Intersection of Eigenvalues. *Proceedings of the 44rd IEEE Conference on Decision and Control December 12-15, 2005, Seville, (Spain). Also on 'Control Systems: Theory, Numerics and Applications, Roma, Italia 30 Mar - 1 Apr 2005, POS, Proceeding of science.*
- [2] A. Agrachev and Y. L. Sachkov. *Control theory from the geometric viewpoint*. Encyclopaedia of Mathematical Sciences, 87. Control Theory and Optimization, II. Springer-Verlag, Berlin, 2004.
- [3] A. Agrachev and A. V. Sarychev. Navier-Stokes equations: controllability by means of low modes forcing. *J. Math. Fluid Mech.*, 7(1):108–152, 2005.
- [4] S.A. Avdonin. On the question of Riesz bases of exponential functions in  $L^2$ . *Vestnik Leningrad. Univ. No. 13 Mat. Meh. Astronom. Vyp. 3, 154*, pages 5–12, 1974.
- [5] J.M. Ball, J.E. Marsden, and M. Slemrod. Controllability for distributed bilinear systems. *SIAM J. Control and Optim.*, 20, July 1982.
- [6] C. Bardos, G. Lebeau, and J. Rauch. Sharp sufficient conditions for the observation, control, stabilization of waves from the boundary. *SIAM J. Cont. Optim.*, 30(5):1024–1065, 1992.
- [7] L. Baudouin. A bilinear optimal control problem applied to a time dependent Hartree-Fock equation coupled with classical nuclear dynamics. *Port. Math. (N.S.)*, 63(3):293–325, 2006.
- [8] L. Baudouin, O. Kavian, and J.-P. Puel. Regularity for a Schrödinger equation with singular potential and application to bilinear optimal control. *J. Diff. Eq.*, 216:188–222, 2005.
- [9] L. Baudouin and J. Salomon. Constructive solutions of a bilinear control problem for a Schrödinger equation. *Systems and Control Letters*, 57(6):453–464, 2008.
- [10] K. Beauchard. Local Controllability of a 1-D Schrödinger equation. *J. Math. Pures et Appl.*, 84:851–956, July 2005.
- [11] K. Beauchard. Controllability of a quantum particule in a 1D variable domain. *ESAIM:COCV*, 14(1):105–147, 2008.
- [12] K. Beauchard. Local Controllability of a 1-D beam equation. *SIAM J. Control Optim.*, 47(3):1219–1273, 2008.

- [13] K. Beauchard. Local Controllability of a 1-D bilinear Schrödinger equation: a simpler proof. (*preprint*), 2009.
- [14] K. Beauchard and J.-M. Coron. Controllability of a quantum particle in a moving potential well. *J. Functional Analysis*, 232:328–389, 2006.
- [15] K. Beauchard, J.-M. Coron, and P. Rouchon. Controllability issues for continuous-spectrum systems and ensemble controllability of Bloch equations. (*preprint*), *arXiv:0903.2720v1*, 2009.
- [16] K. Beauchard and M. Mirrahimi. Practical stabilization of a quantum particle in a one-dimensional infinite square potential well. *SIAM J. Contr. Optim.*, 48(2):1179–1205, 2009.
- [17] N. Burq. Contrôlabilité exacte des ondes dans des ouverts peu réguliers. *Asymptot. Anal.*, 14(2):157–191, 1997.
- [18] N. Burq and P. Gérard. Condition nécessaire et suffisante pour la contrôlabilité exacte des ondes. *C. R. Acad. Sci. Sér. I Math.*, 325(7):749–752, 1997.
- [19] T. Chambrion, P. Mason, M. Sigalotti, and M. Boscain. Controllability of the discrete-spectrum Schrödinger equation driven by an external field. *Ann. Inst. H. Poincaré Anal. Non Linéaire*, 26(1):329–349, 2009.
- [20] J.-M. Coron. On the small-time local controllability of a quantum particle in a moving one-dimensional infinite square potential well. *C. R. Acad. Sciences Paris, Ser. I*, 342:103–108, 2006.
- [21] J.-M. Coron. *Control and nonlinearity*, volume 136. Mathematical Surveys and Monographs, 2007.
- [22] D’Alessandro. *Introduction to quantum control and dynamics*. Applied Mathematics and Nonlinear Science, Series. Boca Raton, FL:Chapman, Hall/CRC, 2008.
- [23] E. Cancès and C. Le Bris and M. Pilot. Contrôle optimal bilinéaire d’une équation de Schrödinger. *CRAS Paris*, 330:567–571, 2000.
- [24] S. Ervedoza and J.-P. Puel. Approximate controllability for a system of schrödinger equations modeling a single trapped ion,. *Annales de l’Institut Henri Poincaré : Analyse non linéaire (to appear)*, 2009.
- [25] A. Fursikov and O. Y. Imanuvilov. Controllability of evolution equations. *Lecture Notes Series, Seoul National University Research Institute of Mathematics Global Analysis Research Center, Seoul*, 34, 1996.
- [26] M. Horvath and I. Joo. On Riesz bases II. *Ann. Univ. Sci. Budapest. Eotvos Sect. Math.*, 33:267–271, 1990.
- [27] R. Ilner, H. Lange, and H. Teismann. Limitations on the control of Schrödinger equations. *ESAIM:COCV*, 12(4):615–635, 2006.
- [28] A. Y. Khapalov. Bilinear controllability properties of a vibrating string with variable axial load and damping gain. *Dyn. Contin. Impuls. Syst. Ser A Math Anal.*, 10(5):721–743, 2003.
- [29] A. Y. Khapalov. Controllability of the semilinear parabolic equation governed by a multiplicative control in the reaction term: a qualitative approach. *SIAM J. Contr. Optim.*, 41(6):1886–1900, 2003.

- [30] A. Y. Khapalov. Controllability properties of a vibrating string with variable axial load. *Discrete Contin. Dyn. Syst.*, 11(2-3):311–324, 2004.
- [31] A. Y. Khapalov. Reachability of nonnegative equilibrium states for the semilinear vibrating string by varying its axial load and the gain of damping. *ESAIM:COCV*, 12:231–252, april 2006.
- [32] A. Y. Khapalov. Local controllability for a ‘swimming’ model. *SIAM J. Contr. Optim.*, 46(2):655–682, 2007.
- [33] V. Komornik. *Exact controllability and stabilization*. Research in Applied Mathematics, Masson, Paris, 1994.
- [34] J. Louis Lions. *Contrôlabilité exacte, perturbations et stabilisation de systèmes distribués, Tome 1*, volume 8. Recherche en Mathématiques Pures et Appliquées [Research in Pure and Applied mathematics], Masson, Paris, 1988.
- [35] M. Mirrahimi. Lyapunov control of a quantum particle in a decaying potential. *Ann. Inst. H. Poincaré (c) Nonlinear Analysis*, 26:1743–1765, 2009.
- [36] M. Mirrahimi and P. Rouchon. Controllability of quantum harmonic oscillators. *IEEE Trans. Automatic Control*, 49(5):745–747, 2004.
- [37] V. Nersesyan. Global approximate controllability for Schrödinger equation in higher Sobolev norms and applications. (*preprint*), 2009.
- [38] V. Nersesyan. Growth of Sobolev norms and controllability of Schrödinger equation. *Comm. Math. Phys. (to appear)*, 2009.
- [39] D. L. Russell. Controllability and stabilisation theory for linear partial differential equations: recent progress and open questions. *SIAM Rev.*, 20(4):639–739, 1978.
- [40] A. Shirikyan. Approximate controllability of three-dimensional Navier-Stokes equations. *Comm. Math. Phys.*, 266(1):123–151, 2006.
- [41] G. Turinici. On the controllability of bilinear quantum systems. In *C. Le Bris and M. Defranceschi, editors, Mathematical Models and Methods for Ab Initio Quantum Chemistry*, volume 74 of Lecture Notes in Chemistry, Springer, 2000.
- [42] E. Zuazua. Exact controllability for semilinear wave equations in one space dimension. *Ann. Inst. H. Poincaré Anal. Non Linéaire*, 10(1):109–129, 1993.